

# The Thermodynamic Web

Chapter-5

# Types of Thermodynamic Properties

- Measured Properties
  - ✓ *P, v, T, Composition*
- Fundamental Properties
  - ✓ *u, s which are the basis for the 1st and 2nd laws of thermodynamics*
- Derived Thermodynamic Properties
  - ✓ *A combination of the above two types of properties,*
    - *H=u+Pv*
    - *a=u-Ts (Helmholz energy)*
    - *g=h-Ts (Gibbs energy)*

# Thermodynamic Property Relationships

- Objective: To develop mathematical expressions through which we can relate the properties of a system to one another and to forms in which data are reported.
- State Postulate:

Knowing two independent intensive properties ( $x,y$ ) constrains the state of a system of constant composition. Then, change in any intensive property such as  $z$  is:

$$dz = \left(\frac{\partial z}{\partial x}\right)_y dx + \left(\frac{\partial z}{\partial y}\right)_x dy \quad du = \left(\frac{\partial u}{\partial T}\right)_v dT + \left(\frac{\partial u}{\partial v}\right)_T dv \quad u=u(T,v)$$



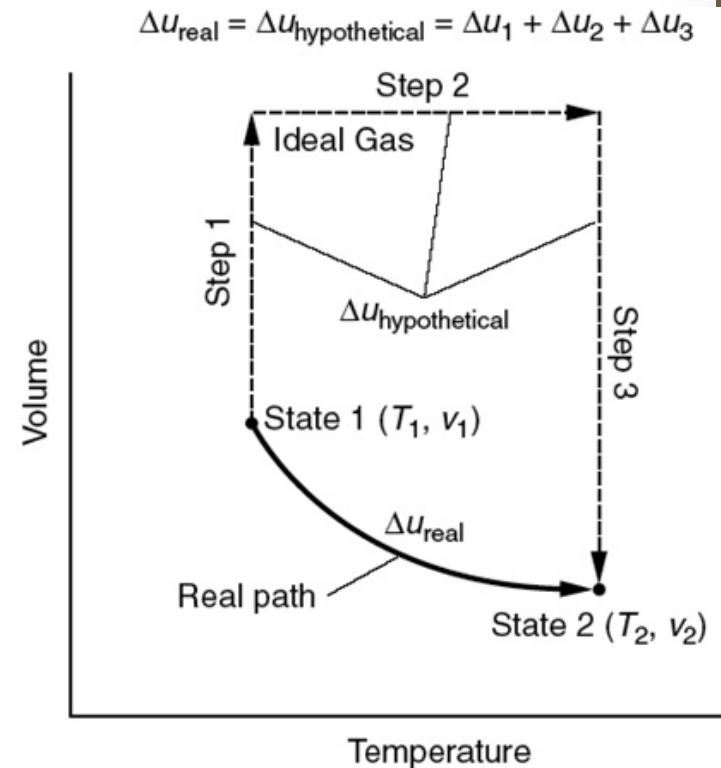
change in the dependent variable  $z$  with respect to the independent variable  $x$  over a path where independent  $y$  is constant

# Thermodynamic Property Relationships

Question: Why should we derive the following equation? useful?

$$du = \left(\frac{\partial u}{\partial T}\right)_v dT + \left(\frac{\partial u}{\partial v}\right)_T dv$$

- We want to calculate  $\Delta u$  for a nonideal gas.
  - $u$  is a state function, path-independent.
  - We can choose hypothetical paths.
  - We only have  $C_p$ ,  $C_v$  data for ideal gases.



# Fundamental Property Relations

- Fundamental Property Relations:

$$du = \delta q_{\text{rev}} + \delta w_{\text{rev}} = T ds - P dv$$

$$dh = du + d(Pv) = T ds + v dP$$

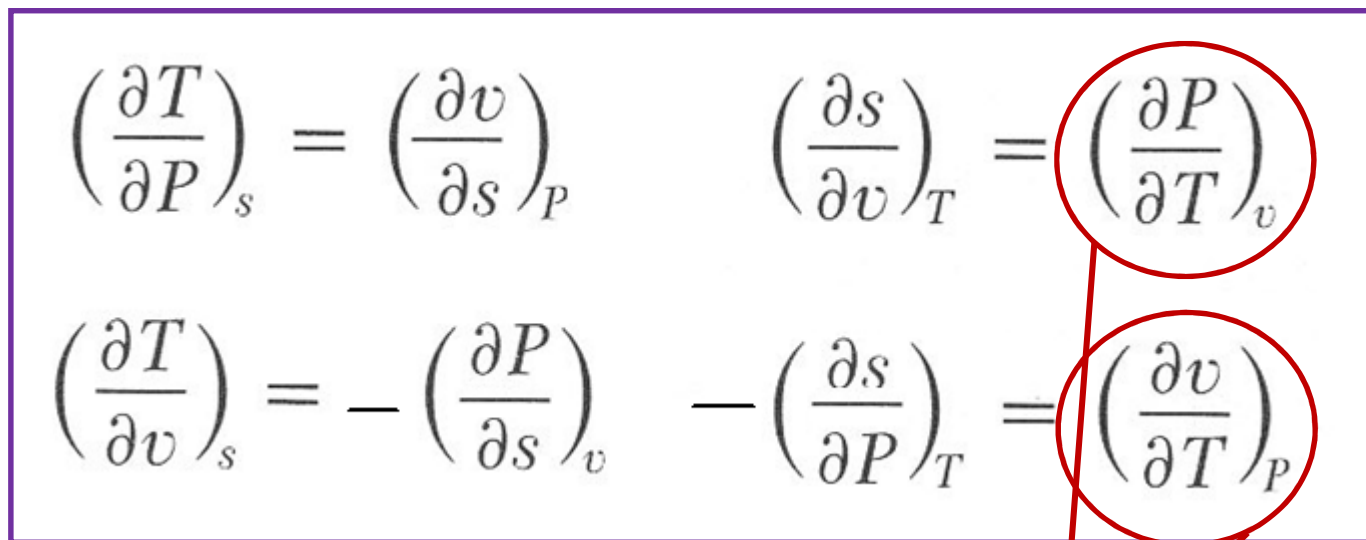
$$da = du - d(Ts) = -s dT - P dv$$

$$dg = dh - d(Ts) = -s dT + v dP$$

$$du = \left(\frac{\partial u}{\partial s}\right)_v ds + \left(\frac{\partial u}{\partial v}\right)_s dv \longrightarrow \left(\frac{\partial u}{\partial s}\right)_v = T \quad \text{and} \quad \left(\frac{\partial u}{\partial v}\right)_s = -P$$

# Maxwell Relations

The **second derivative** of the **fundamental property relationships** are referred to as the Maxwell Relationships.

$$\left(\frac{\partial T}{\partial P}\right)_s = \left(\frac{\partial v}{\partial s}\right)_P \quad \left(\frac{\partial s}{\partial v}\right)_T = \left(\frac{\partial P}{\partial T}\right)_v$$
$$\left(\frac{\partial T}{\partial v}\right)_s = -\left(\frac{\partial P}{\partial s}\right)_v \quad -\left(\frac{\partial s}{\partial P}\right)_T = \left(\frac{\partial v}{\partial T}\right)_P$$
The diagram shows four Maxwell relations arranged in two rows. The first row contains two equations:  $\left(\frac{\partial T}{\partial P}\right)_s = \left(\frac{\partial v}{\partial s}\right)_P$  and  $\left(\frac{\partial s}{\partial v}\right)_T = \left(\frac{\partial P}{\partial T}\right)_v$ . The second row contains two equations:  $\left(\frac{\partial T}{\partial v}\right)_s = -\left(\frac{\partial P}{\partial s}\right)_v$  and  $-\left(\frac{\partial s}{\partial P}\right)_T = \left(\frac{\partial v}{\partial T}\right)_P$ . The terms  $\left(\frac{\partial P}{\partial T}\right)_v$  and  $\left(\frac{\partial v}{\partial T}\right)_P$  are circled in red. Red arrows point from these two circled terms down to a yellow box labeled "Measured Properties".

Measured Properties

# Thermodynamic Web

Problem solving in thermodynamics  $\rightarrow$  Hypothetical Paths  $\rightarrow$   
 to find the change in a property between two states  $\rightarrow$  Partial  
 derivative of a property with respect to another.

**Objective:** Use thermodynamic web to translate partial derivatives  
 as  $C_v$ ,  $C_p$ ,  $\beta$ , and  $\kappa$ .

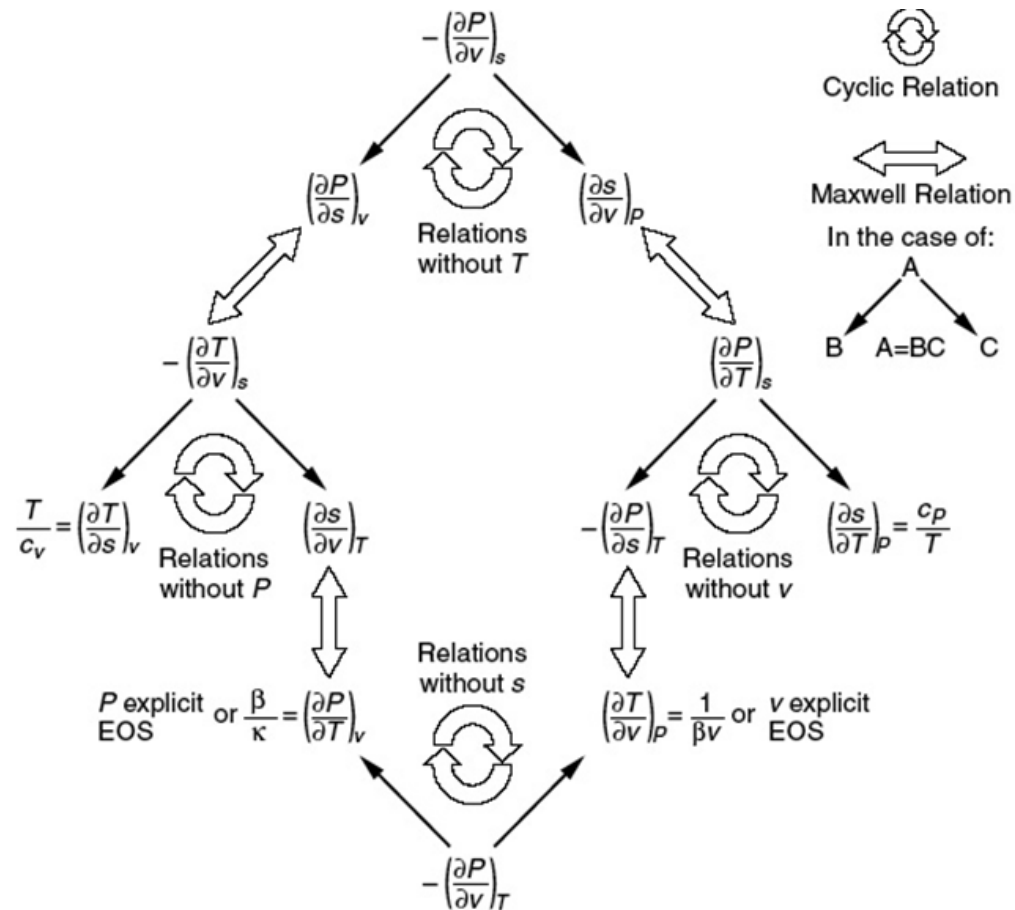
$$-\left(\frac{\partial P}{\partial V}\right)_s = \left(\frac{\partial P}{\partial S}\right)_v \times \left(\frac{\partial S}{\partial V}\right)_P$$

$$-\left(\frac{\partial T}{\partial V}\right)_s \quad \left(\frac{\partial P}{\partial T}\right)_s$$

$$\frac{T}{C_v} \times \left(\frac{\partial S}{\partial V}\right)_T \quad -\left(\frac{\partial P}{\partial S}\right)_T \times \frac{C_p}{T}$$

$$\frac{\beta}{\kappa} \quad \frac{1}{\beta v}$$

$$-\left(\frac{C_p}{v\kappa C_v}\right)$$



$S$  in terms of measured properties,  $v$  and  $T$ :

➤ Mathematics:

$$dz = \left(\frac{\partial z}{\partial x}\right)_y dx + \left(\frac{\partial z}{\partial y}\right)_x dy$$

$$ds = \left(\frac{\partial s}{\partial T}\right)_v dT + \left(\frac{\partial s}{\partial v}\right)_T dv \quad \longrightarrow \quad ds = \frac{c_v}{T} dT + \left(\frac{\partial P}{\partial T}\right)_v dv$$

$$\left(\frac{\partial s}{\partial T}\right)_v = \frac{c_v}{T} \quad \left(\frac{\partial P}{\partial T}\right)_v dv$$

**Maxwell Relation**

$u$  in terms of measured properties,  $v$  and  $T$ :

$$du = \left(\frac{\partial u}{\partial T}\right)_v dT + \left(\frac{\partial u}{\partial v}\right)_T dv \quad \left(\frac{\partial u}{\partial v}\right)_T = \left[ T \left(\frac{\partial P}{\partial T}\right)_v - P \right]$$

$$\left(\frac{\partial u}{\partial T}\right)_v = c_v \quad \left(\frac{\partial u}{\partial v}\right)_T = \left(\frac{T\partial s - P\partial v}{\partial v}\right)_T = \left[ T \left(\frac{\partial s}{\partial v}\right)_T - P \right]$$

$$\longrightarrow du(T, v) = c_v dT + \left[ T \left(\frac{\partial P}{\partial T}\right)_v - P \right] dv$$

# $dh$ in terms of $T$ and $P$

$$dh = \left( \frac{\partial h}{\partial T} \right)_P dT + \left( \frac{\partial h}{\partial P} \right)_T dP \quad \text{and}$$

$$\left( \frac{\partial h}{\partial T} \right)_P = c_P$$

Maxwell

$$h = Tds + vdP \quad \Rightarrow \quad \left( \frac{\partial h}{\partial P} \right)_T = \left( \frac{T\partial s + v\partial P}{\partial P} \right)_T = T \left( \frac{\partial s}{\partial P} \right)_T + v = -T \left( \frac{\partial v}{\partial T} \right)_P + v$$

$$\Rightarrow dh = c_P dT + \left[ -T \left( \frac{\partial v}{\partial T} \right)_P + v \right] dP$$

# Example 5.2

One mole of propane gas is to be expanded from  $0.001 \text{ m}^3$  to  $0.040 \text{ m}^3$  while in contact with a heating bath at  $100^\circ\text{C}$ . The expansion is not reversible. The heat extracted from the bath is  $600 \text{ J}$ . Using the van der Waals equation of state, determine the work for the expansion.

$$\Delta u = q + w$$

$\downarrow$   
 $600 \text{ J}$

$$du = \left(\frac{\partial u}{\partial T}\right)_v dT + \left(\frac{\partial u}{\partial v}\right)_T dv = c_v dT + \left[ T \left(\frac{\partial P}{\partial T}\right)_v - P \right] dv$$

Thermo. Web

$$\Delta u = \int \left[ T \left(\frac{\partial P}{\partial T}\right)_v - P \right] dv$$

$$P = \frac{RT}{v-b} - \frac{a}{v^2} \longrightarrow \left(\frac{\partial P}{\partial T}\right)_v = \frac{R}{v-b}$$

$$\Delta u = \int \left[ \frac{a}{v^2} \right] dv$$

READ EXAMPLE 5.3

# Joule-Thomson Expansion

➤ Unrestrained, free expansion of real gases.

□ Analysis:

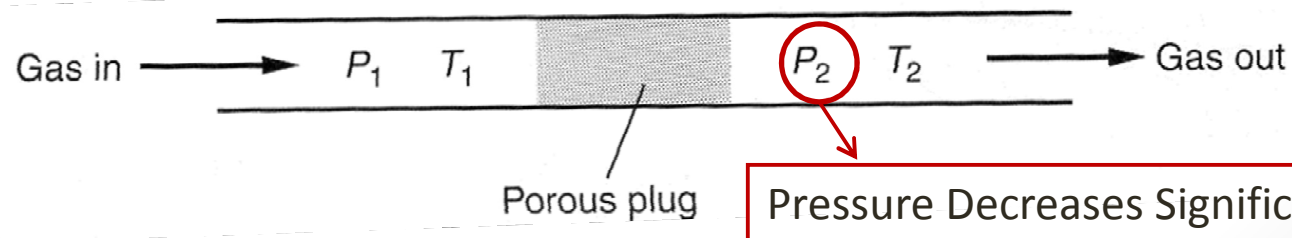
- ✓ Gas spends little time in the plug, thus no heat transfer i.e. adiabatic.
- ✓ Shaft work=0

❖ *First Law at steady state:*

$$H_2 - h_1 = \Delta H = 0 \quad \text{isenthalpic}$$

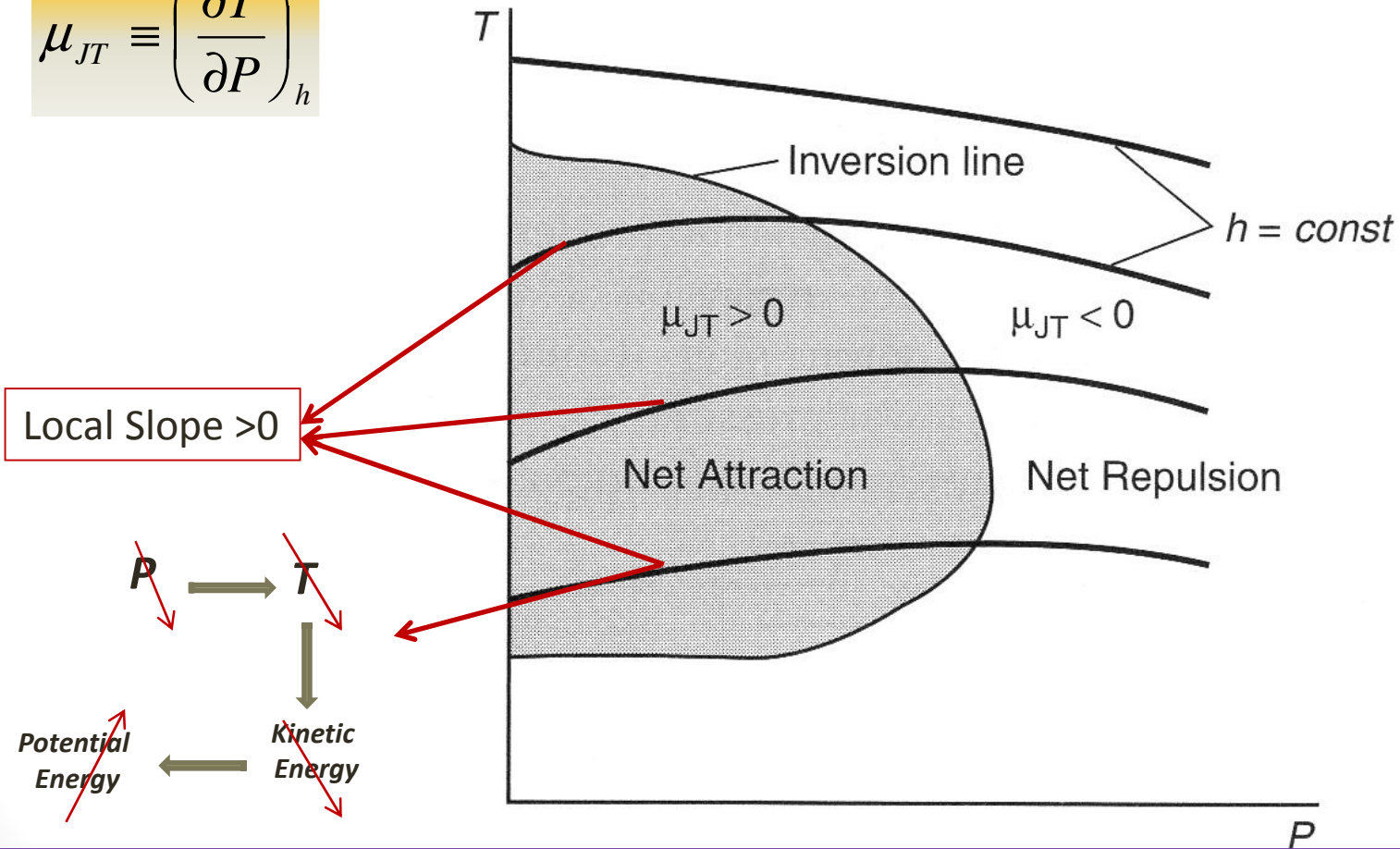
We can determine the change in temperature as the pressure decreases in an isenthalpic process if we know the Joule-Thomson coefficient:

$$\mu_{JT} \equiv \left( \frac{\partial T}{\partial P} \right)_h$$



# Joule-Thomson Expansion

$$\mu_{JT} \equiv \left( \frac{\partial T}{\partial P} \right)_h$$



**Boyle Temperature:** For a given pressure, the temperature at which attractive and repulsive interactions balance.

# Expression for $\mu_{JT}$

$$dh = c_P dT + \left[ -T \left( \frac{\partial v}{\partial T} \right)_P + v \right] dP$$

$$c_P^{\text{real}} = c_P(T, P) = c_P^{\text{ideal gas}} - \int_{P_{\text{ideal gas}}}^{P_{\text{real}}} \left[ T \left( \frac{\partial^2 v}{\partial T^2} \right)_P \right] dP$$

since interactions  
are important

$$\mu_{JT} = \left( \frac{\partial T}{\partial P} \right)_h = \frac{\left[ T \left( \frac{\partial v}{\partial T} \right)_P - v \right]}{c_P} = \frac{\left[ T \left( \frac{\partial v}{\partial T} \right)_P - v \right]}{c_P^{\text{ideal gas}} - \int_{P_{\text{ideal gas}}}^{P_{\text{real}}} \left[ T \left( \frac{\partial^2 v}{\partial T^2} \right)_P \right] dP}$$

If we have an equation of state and the heat capacity of a fluid, we can evaluate  $\mu_{JT}$ .

# Example 5.7

Develop an expression for the Joule–Thomson coefficient using the pressure-based expansion of the virial equation truncated at the second virial coefficient. Use the corresponding state relationships presented in Chapter 4 for the temperature dependence of  $B$  to develop a generalized correlation for  $\mu_{JT}$ .

$$\mu_{JT} = \left(\frac{\partial T}{\partial P}\right)_h = \frac{\left[ T \left(\frac{\partial v}{\partial T}\right)_P - v \right]}{c_P} = \frac{\left[ T \left(\frac{\partial v}{\partial T}\right)_P - v \right]}{c_P^{\text{ideal gas}} - \int_{P_{\text{ideal gas}}}^{P_{\text{real}}} \left[ T \left(\frac{\partial^2 v}{\partial T^2}\right)_P \right] dP}$$

**Virial Eq.:**

$$v = \frac{RT}{P}(1 + B'P) = \frac{RT}{P} + B'RT \rightarrow \begin{cases} \left(\frac{\partial v}{\partial T}\right)_P = \frac{R}{P} + RB' + RT \left(\frac{dB'}{dT}\right) \\ \left(\frac{\partial^2 v}{\partial T^2}\right)_P = R \left(\frac{dB'}{dT}\right) + RT \left(\frac{d^2B'}{dT^2}\right) \end{cases}$$

$$\mu_{JT} = \left(\frac{\partial T}{\partial P}\right)_h = \frac{RT^2 \left(\frac{dB'}{dT}\right)}{c_P^{\text{ideal gas}} - \int_{P_{\text{ideal gas}}}^{P_{\text{real}}} \left\{ T \left[ R \left(\frac{dB'}{dT}\right) + RT \left(\frac{d^2B'}{dT^2}\right) \right] \right\} dP}$$

**Corresponding State Relationship:**

$$B_r = \frac{BP_c}{RT_c} = B^{(0)} + \omega B^{(1)}$$

$$\left\{ \begin{array}{l} B^{(0)} = 0.083 - \frac{0.422}{T_r^{1.6}} \\ B^{(1)} = 0.139 - \frac{0.172}{T_r^{4.2}} \end{array} \right.$$

$$B' = \frac{B}{RT} = \frac{B_r T_c}{P_c T} = \frac{B_r}{P_c T_r} = \frac{B^{(0)} + \omega B^{(1)}}{P_c T_r} \rightarrow B' = \frac{1}{P_c} \left[ \left( \frac{0.083}{T_r} - \frac{0.422}{T_r^{2.6}} \right) + \omega \left( \frac{0.139}{T_r} - \frac{0.172}{T_r^{5.2}} \right) \right]$$

$$\left( \frac{dB'}{dT} \right) = \frac{1}{P_c} \left[ \left( -\frac{0.083}{TT_r} + \frac{1.097}{TT_r^{2.6}} \right) + \omega \left( -\frac{0.139}{TT_r} + \frac{0.894}{TT_r^{5.2}} \right) \right]$$

$$\left( \frac{d^2B'}{dT^2} \right) = \frac{1}{P_c} \left[ \left( \frac{0.166}{T^2 T_r} - \frac{3.950}{T^2 T_r^{2.6}} \right) + \omega \left( \frac{0.278}{T^2 T_r} - \frac{5.545}{T^2 T_r^{5.2}} \right) \right]$$

$$\mu_{JT} = \frac{-\frac{T_c}{P_c} \left[ \left( -0.083 + \frac{1.097}{T_r^{1.6}} \right) + \omega \left( -0.139 + \frac{0.994}{T_r^{4.2}} \right) \right]}{\frac{c_P^{\text{ideal gas}}}{R} - P_r \left[ \left( \frac{0.083}{T_r} - \frac{2.853}{T_r^{2.6}} \right) + \omega \left( \frac{0.139}{T_r} - \frac{4.651}{T_r^{5.2}} \right) \right]}$$

If the critical properties and acentric factor are known, the JT coefficient can be calculated at  $T$  and  $P$ .