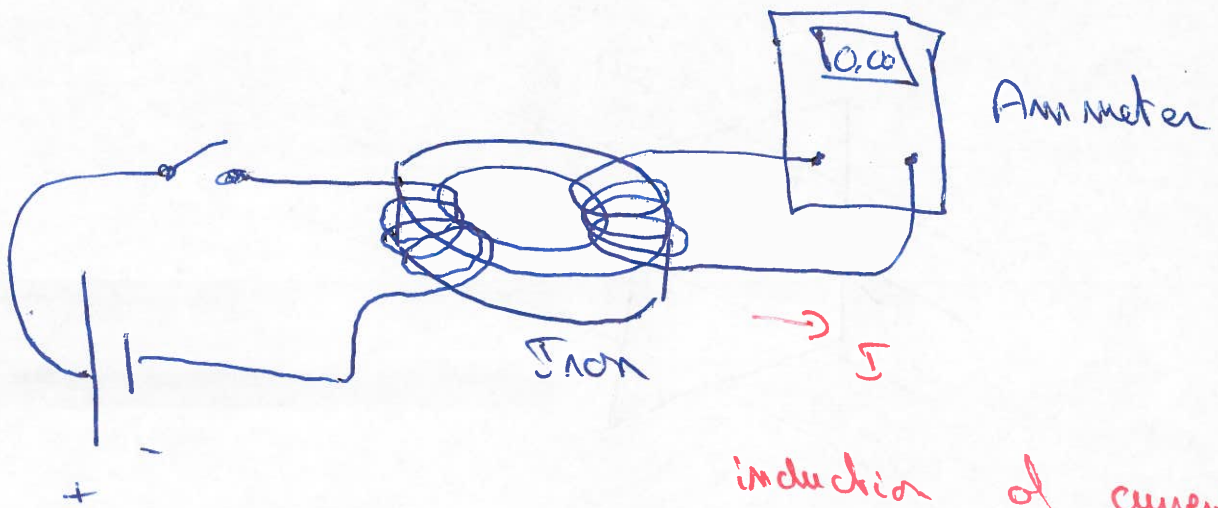


Faraday's law

1

Chapter 31

Faraday law of induction.



- induction of current
- * when switching on the the battery
 - * ————— off the battery
- ⇒ transient

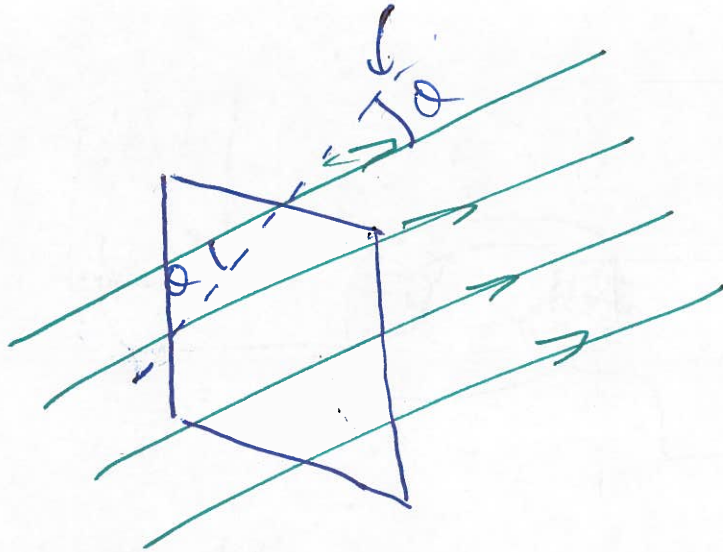
$$\| \mathcal{E} = - \frac{d\Phi_B}{dt} \quad \text{with} \quad \Phi_B = \int \vec{B} \cdot d\vec{S}$$

if a coil consist of N loops , emf in serie

$$\mathcal{E} = -N \frac{d\phi_B}{dt}$$

in a uniform field \vec{B}

normal to the loop

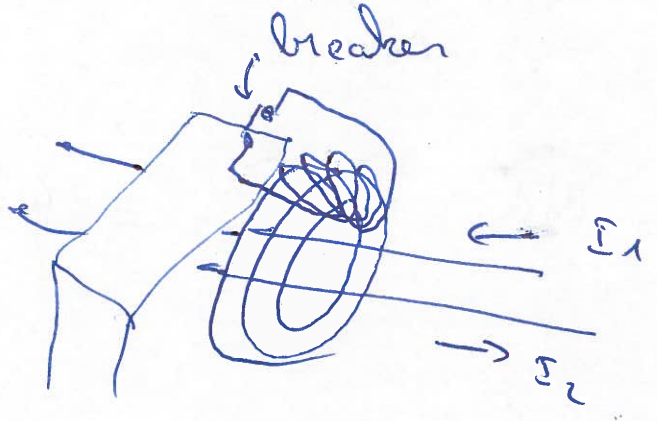


$$\Rightarrow \mathcal{E} = -\frac{d}{dt} (BS \cos \theta)$$

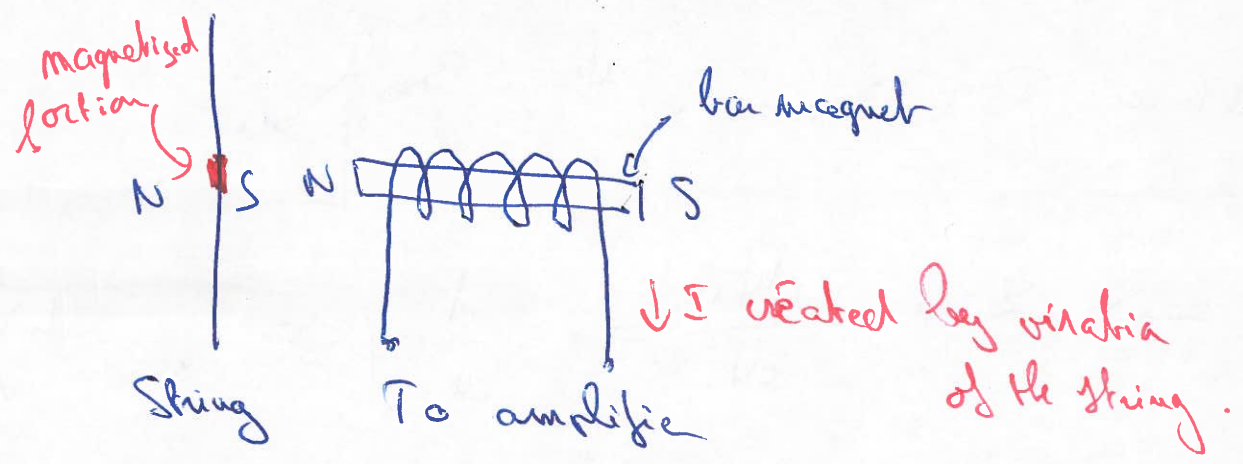
Induction of a current:

- by varying B
- by varying S
- by varying θ

ex: Ground fault circuit interrupter



Guitar pickup coil



Inducing an emf in a coil

coil of 200 loop, each loop is a square ($d = 18cm$)

in a uniform field $\vec{B} \perp$ to the plane of the coil

$B \nearrow$ linearly from 0 to $0,5T$ in $0,80s$

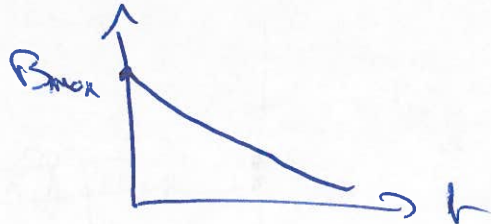
\Rightarrow induced emf ??

$$|\mathcal{E}| = N \frac{\Delta \Phi_B}{\Delta t} = N \frac{\Delta (BS)}{\Delta t}$$

$$|\mathcal{E}| = Nd^2 \frac{B_f - B_i}{\Delta t} = 4.0 \text{ V}$$

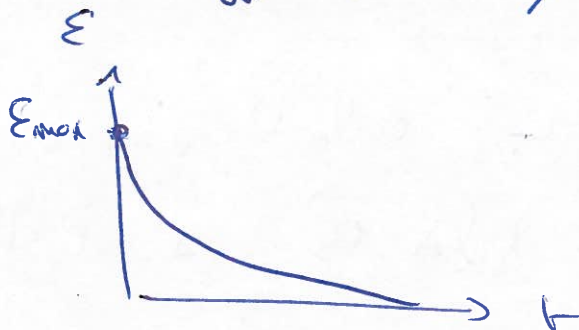
* A loop with a surface S , $\vec{B} \perp$ to the plane of the loop

at $t=0$: $B = B_{\max} e^{-at}$



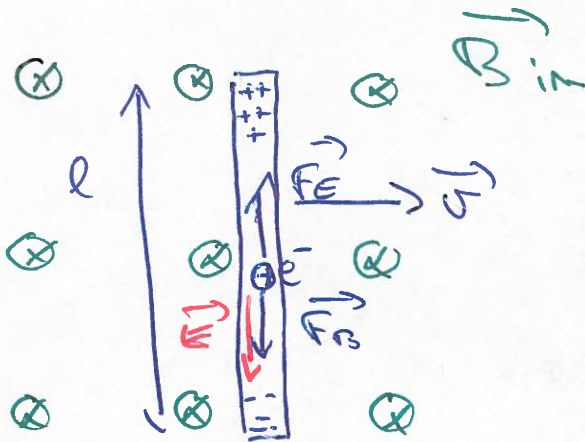
$$\Rightarrow \mathcal{E} = - \frac{d\Phi_B}{dt} = - \frac{d}{dt} (SB_{\max} e^{-at})$$

$$= - SB_{\max} \frac{d}{dt} (e^{-at}) = \underbrace{aS B_{\max}}_{\mathcal{E}_{\max}} e^{-at}$$



Motional emf

3



Straight conductor of length l , moving at a speed \vec{v} because of an external agent

\Rightarrow on each free electron:

$$\vec{F}_B = q\vec{v} \times \vec{B}$$

\Rightarrow generation of an electric field \vec{E} , $\vec{F}_E = q\vec{E}$

Equilibrium $\Rightarrow \vec{F}_B = -\vec{F}_E$

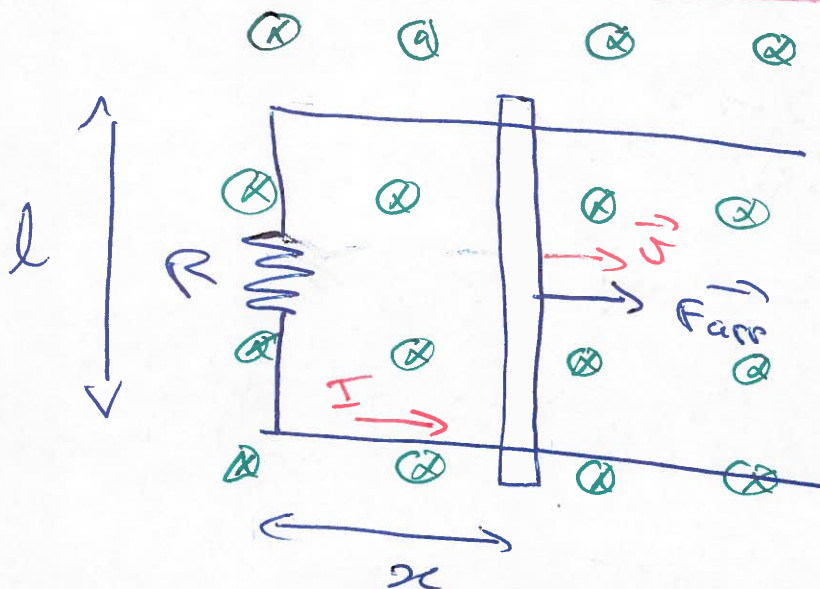
$$\Rightarrow qvB = qE \Rightarrow$$

$$\boxed{E = vB}$$

\vec{E} is uniform $\Rightarrow \Delta V = El$

$$\Rightarrow \boxed{\Delta V = El = Blv}$$

Moving a conductor in a closed conducting path



$$\underline{\underline{\Phi_B = B l x}}$$

$$\Rightarrow \mathcal{E} = - \frac{d\Phi_B}{dt} = - \frac{d}{dt} (B l x) = - B l \frac{dx}{dt}$$

$$\Rightarrow \boxed{\mathcal{E} = - B l v} \quad \begin{array}{l} \text{Motional} \\ \text{emf} \end{array}$$

$$\Rightarrow I = \frac{|\mathcal{E}|}{R} = \frac{B l v}{R}$$

Energy consideration

(4)

A work is applied on the bar

Conservation of energy $W = \Delta E_{int}$

\Rightarrow input energy \Rightarrow internal energy in the resistor

When moving the bar \Rightarrow force \vec{F}_B is experienced

\Rightarrow magnitude ILB

$$\vec{v} = dv/dt, \text{ i.e. } \Rightarrow \underline{\underline{\vec{a} = 0}} \Rightarrow \sum \vec{F} = \vec{0}$$

$$\Rightarrow F_{app} = F_B = ILB$$

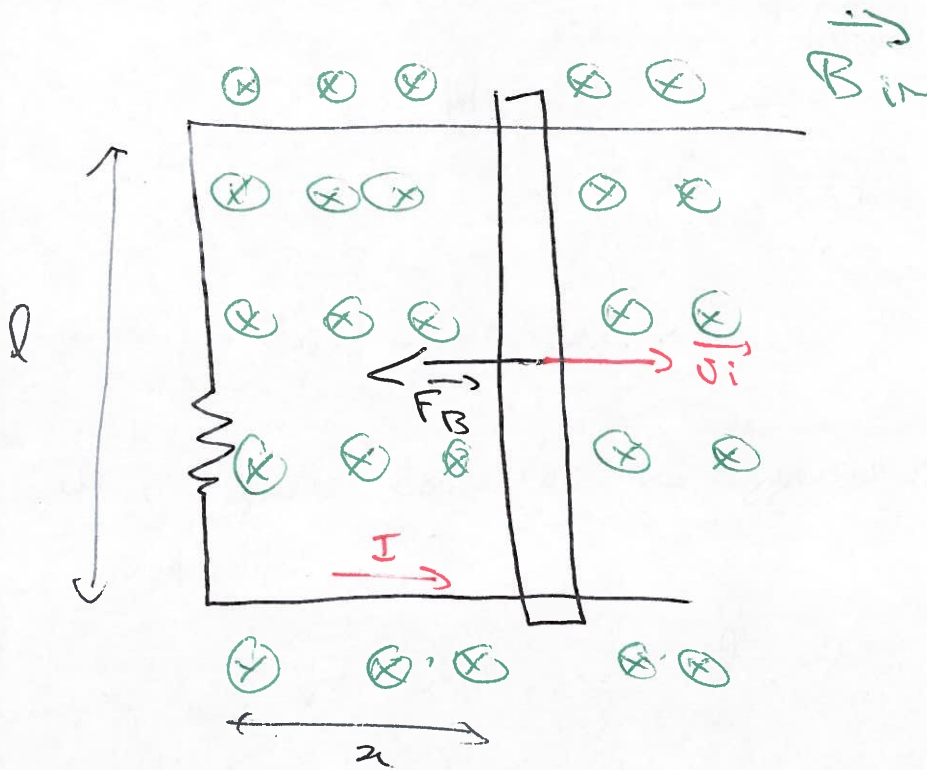
\Rightarrow Power delivered by the applied force

$$P = F_{app} v = (ILB)v = \frac{B^2 l^2 v^2}{R} = \frac{\mathcal{E}^2}{R}$$

$$(W = F_{app} x) \quad \left(I = \frac{Blv}{R} \right)$$

$$P = \frac{dW}{dt}$$

$$\boxed{P = \frac{\mathcal{E}^2}{R}}$$



Initial speed at $t=0$ v_i

$$\vec{F}_B = -I l B \hat{i}$$

$$\sum \vec{F} = m \vec{a} \Rightarrow F_x = m a$$

$$\Rightarrow -I l B = m \frac{dv}{dt}$$

$$I = B l v / R$$

$$\Rightarrow m \frac{dv}{dt} = - \frac{B^2 l^2}{R} v \quad \Rightarrow \quad \frac{dv}{v} = - \frac{B^2 l^2}{m R} dt$$

$$\int_{v_i}^v \frac{dv}{v} = \int_0^t \left(- \frac{B^2 l^2}{m R} \right) dt$$

$$\ln\left(\frac{\sigma}{\sigma_0}\right) = - \frac{B^2 l^2}{m R} t$$

$$\Rightarrow \boxed{\sigma = \sigma_0 e^{-t/\tau} \quad \text{with } \tau = \frac{m R}{B^2 l^2}}$$

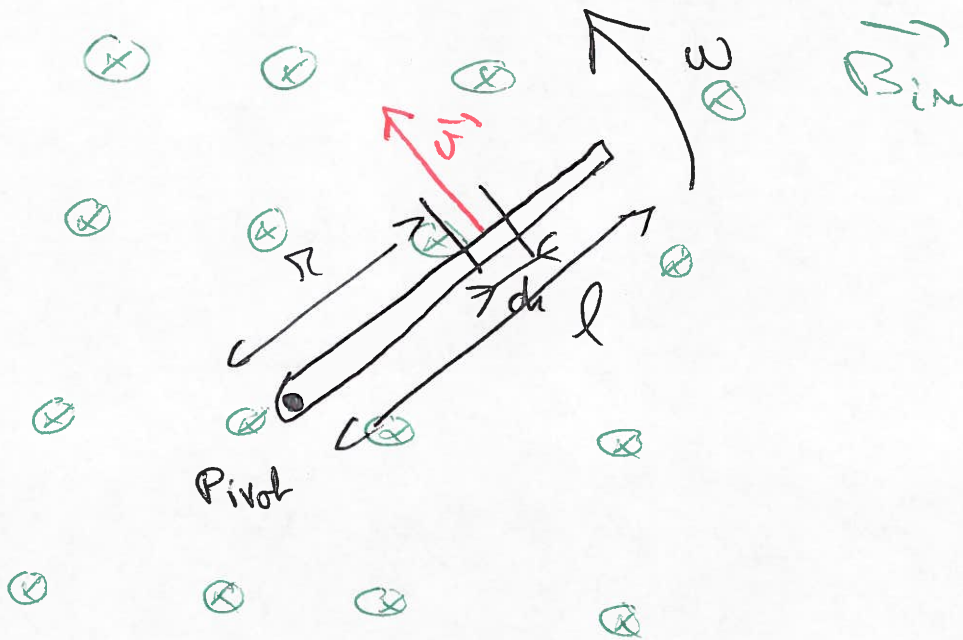
On : $P_{resistor} = - P_{bar}$

$$I^2 R = - \frac{d}{dt} \left(\underbrace{\frac{1}{2} m v^2}_{\Delta K} \right)$$

$$\Rightarrow \frac{B^2 l^2 v^2}{R} = - m v \frac{dv}{dt}$$

$$\Rightarrow \frac{dv}{v} = - \left(\frac{B^2 l^2}{m R} \right) dt$$

Rotational emf induced in a rotating bar



$$|d\varepsilon| = B v dr \quad \left. \vphantom{|d\varepsilon|} \right\} \text{small elements in } \underline{\text{serie}}$$

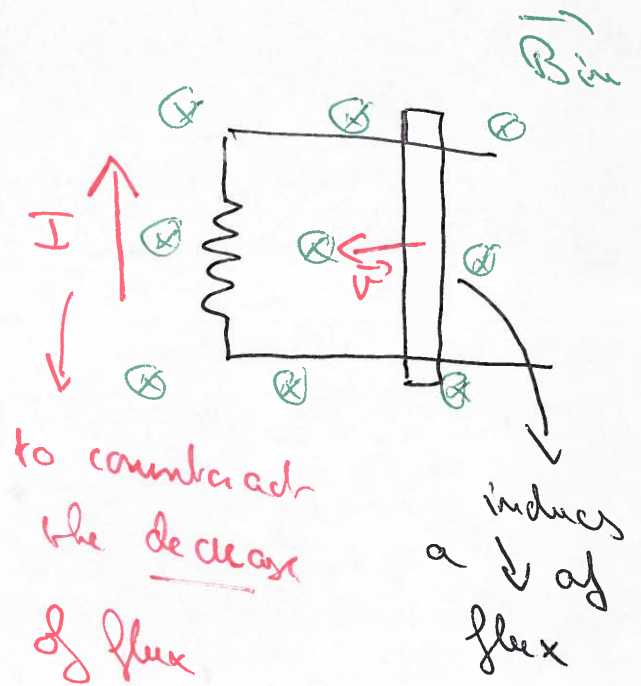
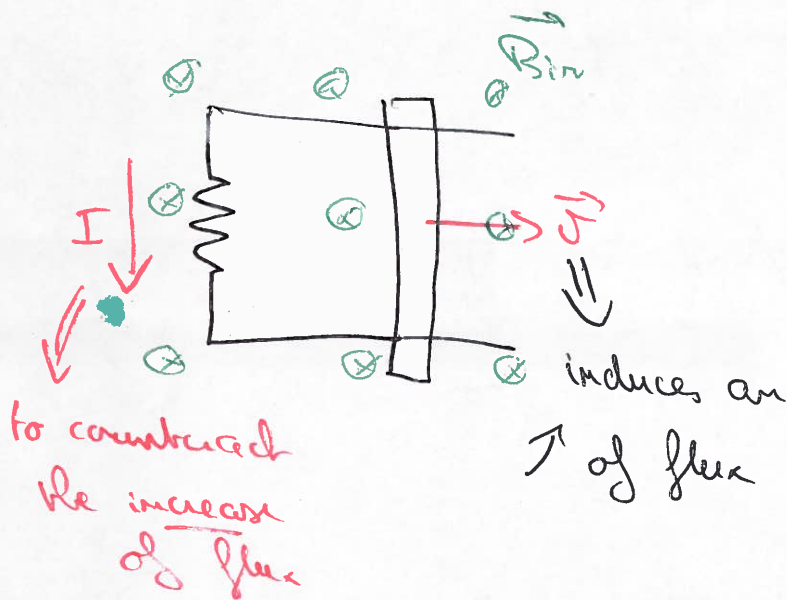
$$\Rightarrow |\varepsilon| = \int B v dr \quad \text{with} \quad v = \omega r$$

$$\Rightarrow |\varepsilon| = B \int v dr = B \omega \int_0^l r dr$$

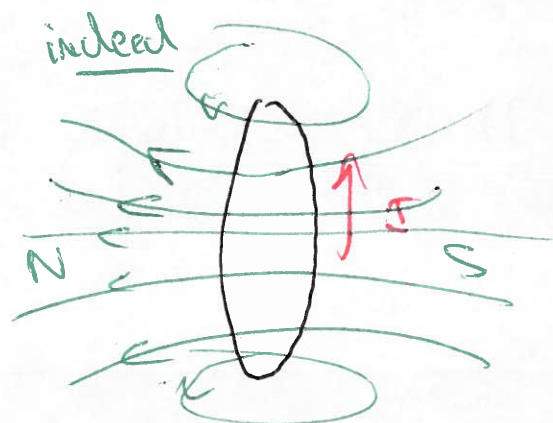
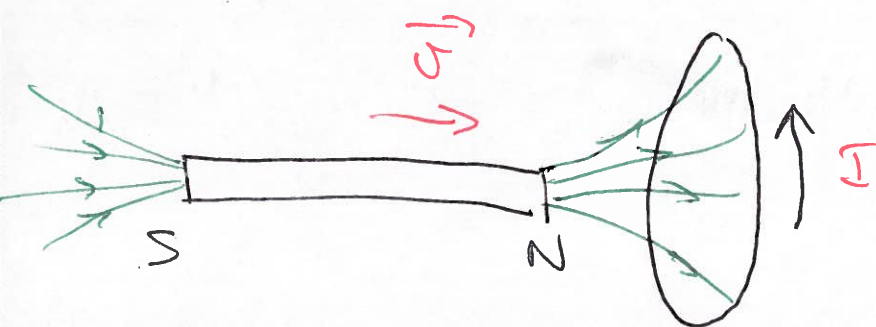
$$|\varepsilon| = \frac{1}{2} B \omega l^2$$

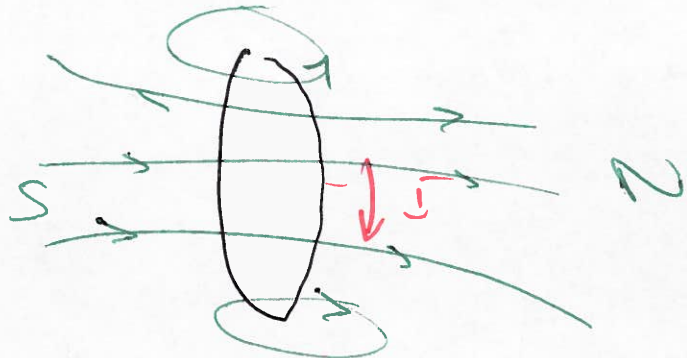
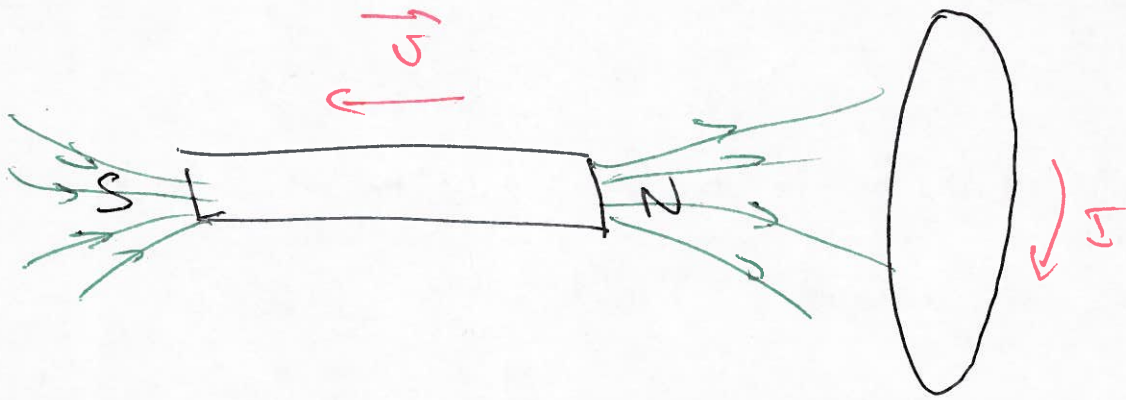
Lenz's Law

The induced current in a loop is in the direction that creates a magnetic field that opposes the change of magnetic flux through the area enclosed by the loop.

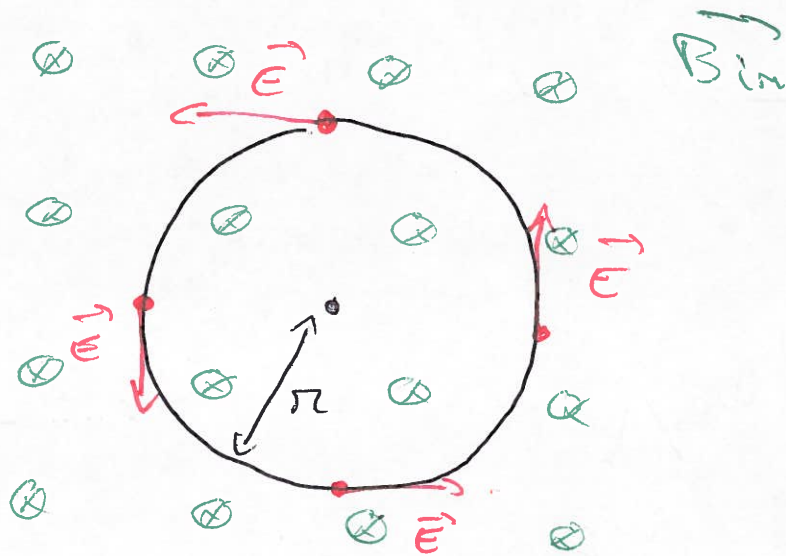


$$\mathcal{E} = - \frac{d\Phi_B}{dt}$$





Induced emf and electric fields



If B is varying with time $\Rightarrow E = -d\Phi_B/dt$

Note: In electrostatic, \vec{E} exists even if there is not net charge \Rightarrow moving charges \vec{E}

\Rightarrow change in magnetic flux \Rightarrow induced \vec{E}

\vec{E} induced is non conservative.

\vec{E} electrostatic (created by static charges) is conservative

Work done by the electric field in moving a charge q around the loop = qE

Force $q\vec{E}$ \Rightarrow work around the loop $qE(2\pi r)$

We should have: $qE = qE(2\pi r)$

$$\Rightarrow E = E / 2\pi r$$

$$E = -d\Phi_B / dt \Rightarrow$$

$$\Phi_B = BS = B\pi r^2$$

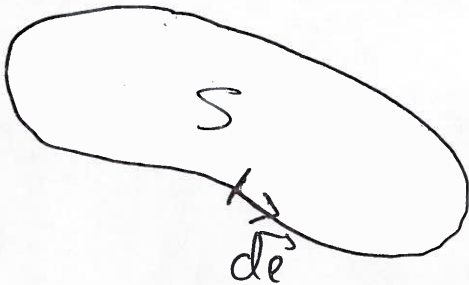
$$E = -\frac{\pi}{2} \frac{dB}{dt}$$

\Rightarrow time varying magnetic field.

\Rightarrow induced electric field

$$\mathcal{E} = \oint \vec{E} \cdot d\vec{l}$$

For any close path \mathcal{C}



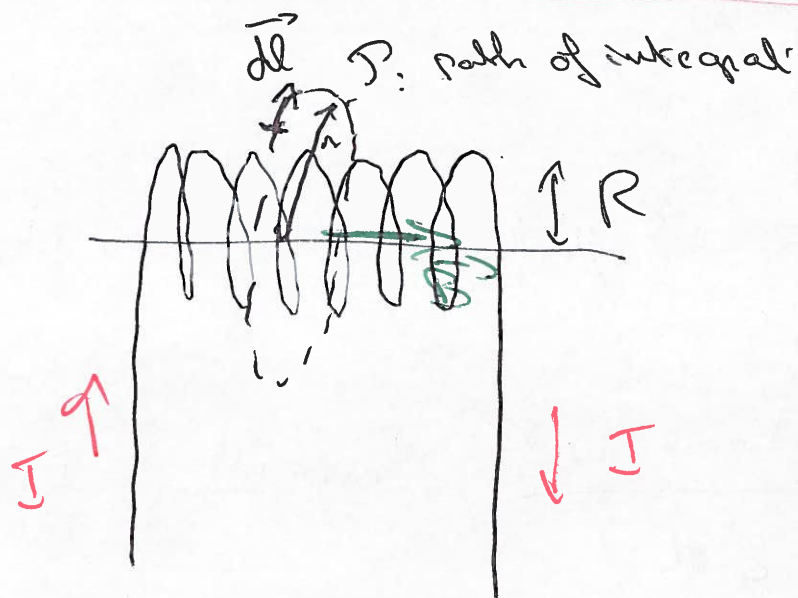
$$\oint_{\mathcal{C}} \vec{E} \cdot d\vec{l} = - \frac{d\phi_B}{dt}$$

Note when \vec{B} is varying \Rightarrow \vec{E} induced is Not conservative.

Note 2 : Electrostatic field : $\Delta V = - \int \vec{E} \cdot d\vec{l}$
 $\Delta V = - \oint \vec{E} \cdot d\vec{l} = \underline{\underline{0}}$
Conservative field

8

Electric field induced by a change
in magnetic field in a solenoid.



\vec{B} uniform inside the solenoid

$$\vec{B} = \mu_0 n I$$

$$\underline{I = I_{\max} \cos \omega t}$$

$$\oint \vec{E} \cdot d\vec{l} = - \frac{d\Phi_B}{dt} = - \frac{d}{dt} (B \pi R^2)$$

$$B = \mu_0 n I, \quad I = I_{\max} \cos \omega t$$

For $\pi > R$

$$\Rightarrow - \frac{d\Phi_B}{dt} = - \pi R^2 \mu_0 n I_{\max} \frac{d}{dt} (\cos \omega t) = \pi R^2 \mu_0 n I_{\max} \omega \sin \omega t$$

$$\oint \vec{E} \cdot d\vec{\ell} = E \int d\ell = E (2\pi r)$$

\downarrow
 constant along
 \int

$$E (2\pi r) = \pi R^2 \mu_0 N I_{\max} \omega \sin \omega t$$

$$E = \frac{\mu_0 N I_{\max} \omega R^2}{2r} \sin \omega t \quad \text{for } r > R$$

$E \propto 1/r$

For $r < R$

$$-\frac{d\Phi_B}{dt} = -\frac{d}{dt} (B \pi r^2)$$

$$E (2\pi r) = \pi r^2 \mu_0 N I_{\max} \omega \sin \omega t$$

$$E = \frac{\mu_0 N I_{\max} \omega}{2} r \sin \omega t \quad \text{for } r < R$$

Generators and Motors

Generator : energy by work \rightarrow electrical transmission

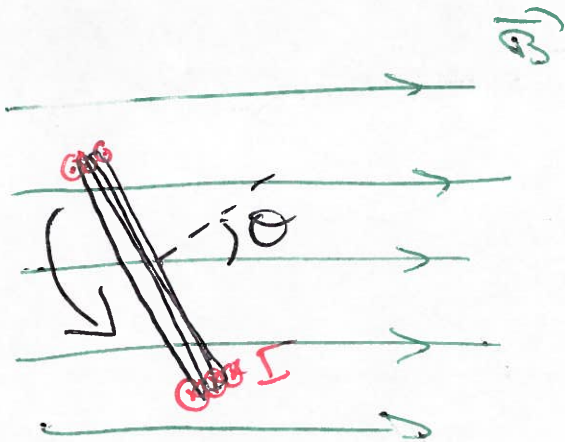
Motor : convert \Rightarrow electrical transmission \Rightarrow
energy transferred out by work.

Generator of AC current

$$\phi_B = B S \cos \theta = B S \cos \omega t$$

N loop

$$\epsilon = - N \frac{d\phi_B}{dt} = NBS \omega \sin \omega t$$

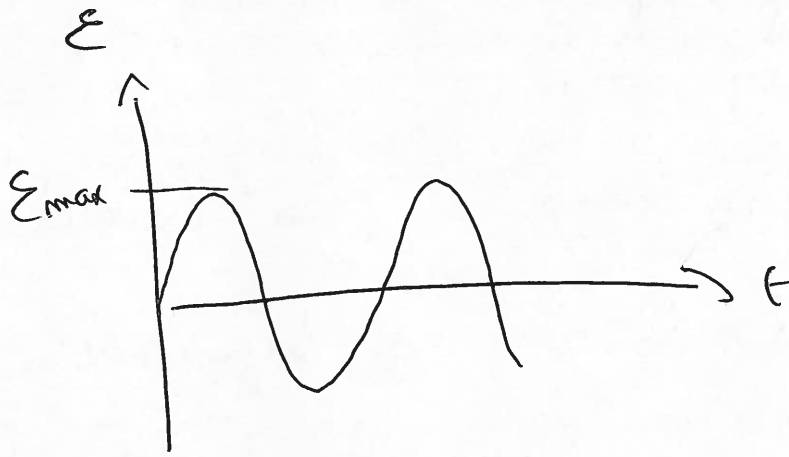


$$\epsilon_{\max} = NBS \omega$$

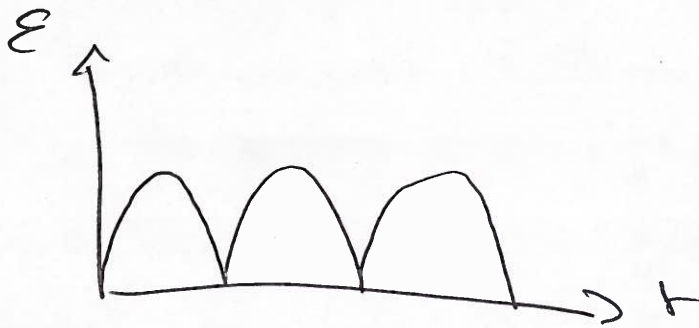
$$\omega = 2\pi f$$

$f = 60 \text{ Hz}$ in USA, Canada

$f = 50 \text{ Hz}$ in Europe



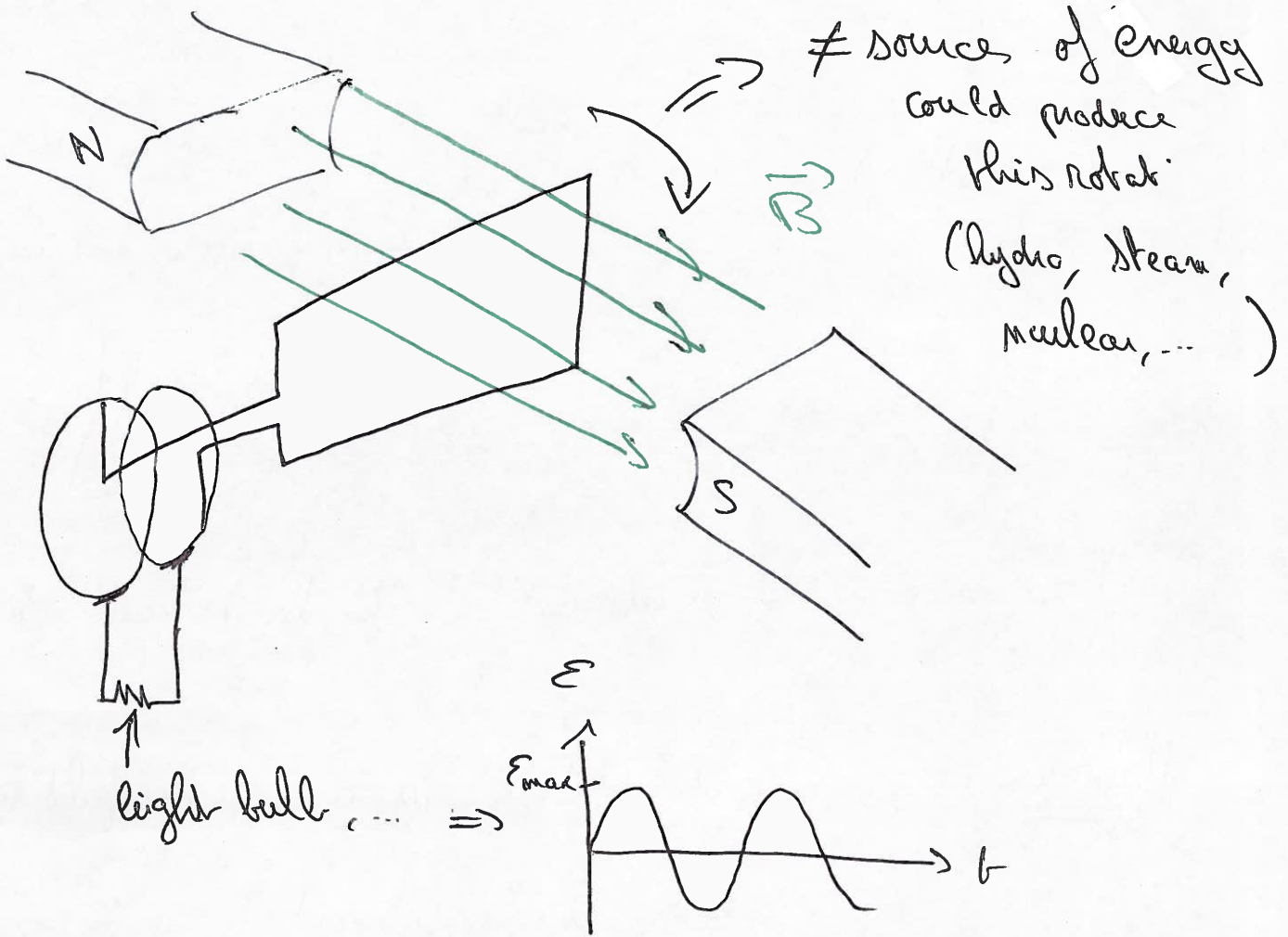
DC current generator : use of a commutator



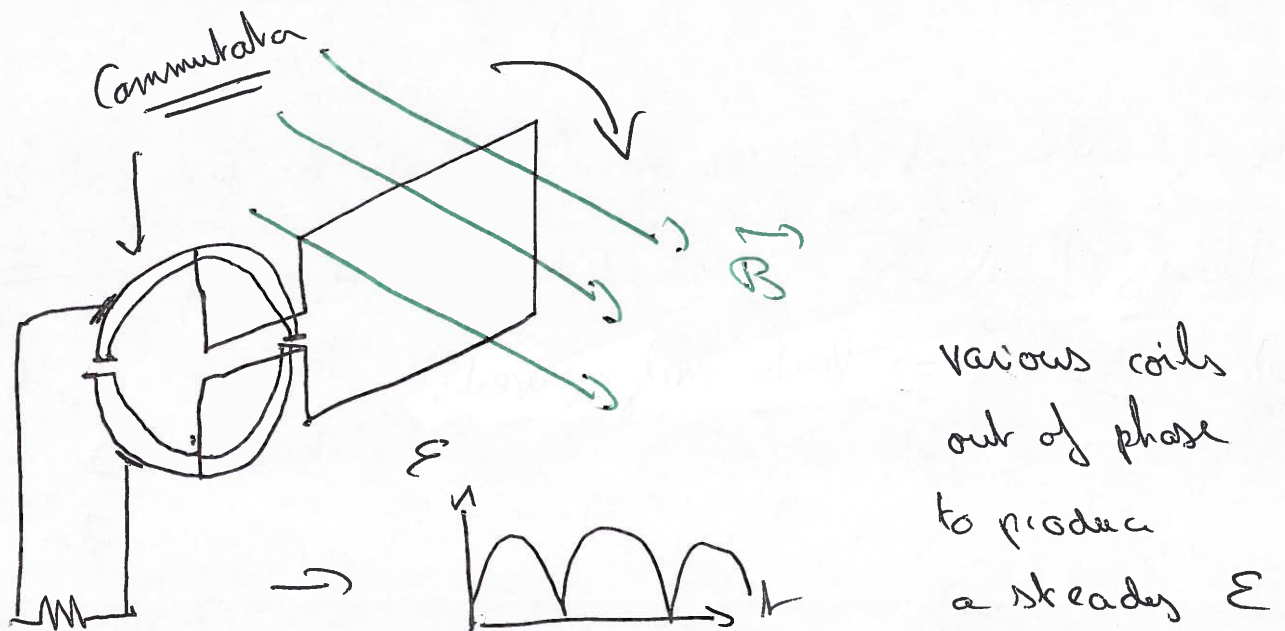
Various coils out of phase to
produce a steady ϵ .

Induced current in a motor.

AC current generator



DC current generator



Motor : electrical energy \Rightarrow transferred to work (rotation)

Current supplied to a coil \Rightarrow torque on the
(from a battery) current carrying coil in \vec{B}

\Rightarrow Rotation

Transmission of mechanical work to external device.

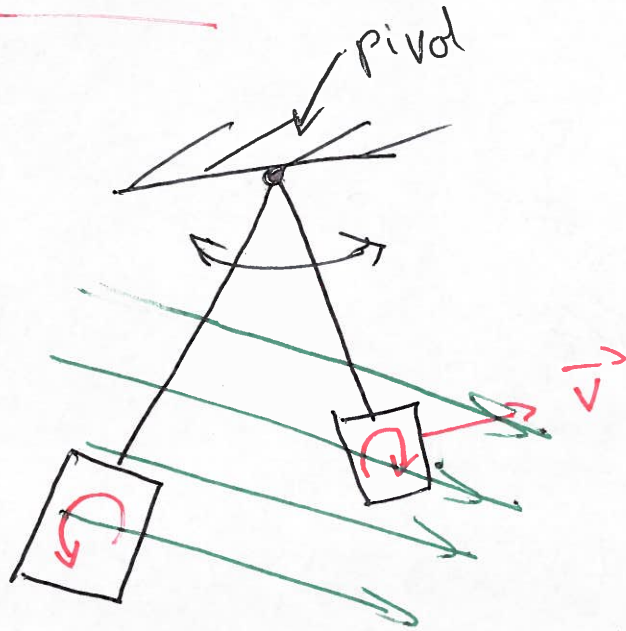
However \Rightarrow coil rotating in a magnetic field \vec{B}

\Rightarrow emf induced

From Lenz's law \Rightarrow back emf to reduce the current in the coil

// If mechanical load $\uparrow \Rightarrow$ motor slows down \Rightarrow back emf \downarrow
 \Rightarrow more current in the coil \Rightarrow more consumption.
Heavy load \Rightarrow motor jam \Rightarrow lack of back emf
can lead to dangerous high current in the motor
 \Rightarrow heat

Eddy currents

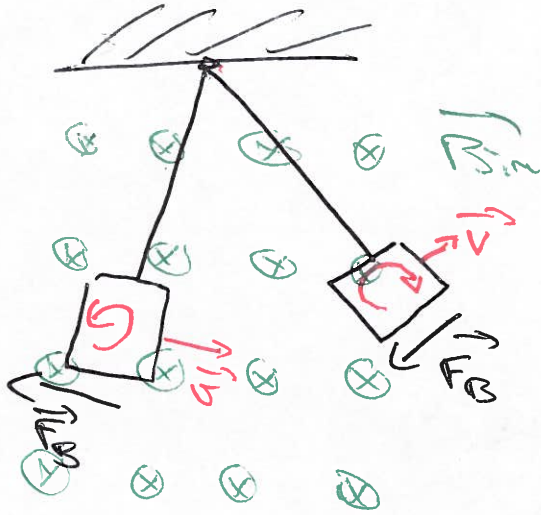


⇒ Movement of piece of metal in a magnetic field ⇒ induced circulating currents

⇒ eddy current

Lenz law ⇒ magnetic field generated by the currents in the plates ⇒ poles that repel by the poles of the magnet ⇒

∥ Repulsive force that oppose to the motion of the plate



- Applications Brakes (train, truck) (brake without contact)
- Gaps in the plates to prevent the creation of eddy currents (to eliminate them since they are undesirable in most cases)