

**ECSE 351B Electromagnetic Fields
Assignment 5**

Due: 22 March 2011

1/2 1. Problem 3-18 from the course text.

1/2 2. Problem 3-19 from the course text.

1 3. Problem 3-24 from the course text.

5 4. Problem 3-26 from the course text.

1 5. Problem 3-28 from the course text.

1 6. A very long (infinite) cylindrical magnet has constant magnetization everywhere inside the magnet equal to $\mathbf{M} = 5000$ (A/m) directed along the axis. The diameter of the magnet is 0.04 (m). Calculate the magnetic flux density due to this magnet everywhere in space. Design an equivalent free-space system (involving only conductors) that produces an identical magnetic flux density everywhere in space.

1
10 7. Consider the flux plot shown below of the static magnetization field \mathbf{M} in a linear material region. Assuming that \mathbf{M} does not vary in the y direction, determine the direction of the net bound volume current density at any given point within the region.

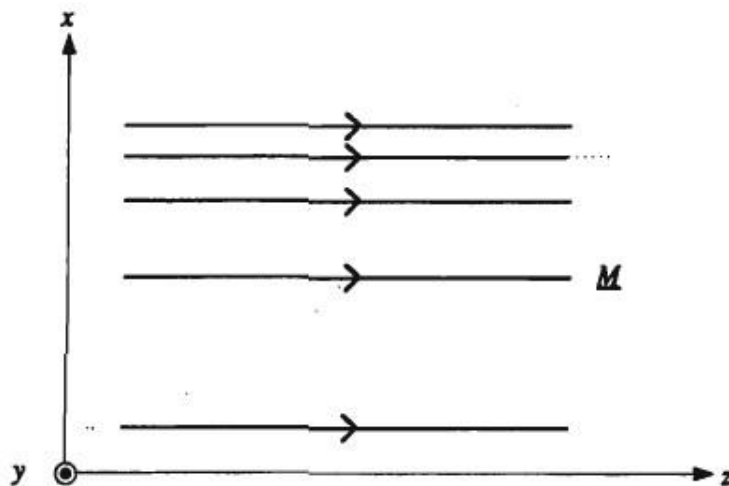
(a) Negative x direction.

(b) Negative y direction.

(c) Positive x direction.

(d) Positive y direction.

(e) I don't know ...



Problems # 1, 2, 3

3-18.) Given $N = 3.3 \times 10^{28}$ atoms/m³,
 $\vec{H} = 4\vec{a}_z$ A/m, $m = \vec{a}_x 1.2 \times 10^{-24}$ A·m².

(10) The volume magnetization density becomes, from (3-55):

$$\vec{M} = N\vec{m} = (3.3 \times 10^{28})\vec{a}_x 1.2 \times 10^{-24}$$

$$= 39.6 \text{ kA/m} \quad \dots(1)$$

making, from (3-60):

$$\chi_m = \frac{M}{H} = \frac{39600}{4} = 9900 \quad \dots(2)$$

so the relative permeability, from (3-62), is

$$\mu_r = 1 + \chi_m = 9901 \quad \dots(3)$$

(b) From (3-64c):

$$\vec{B} = \mu\vec{H} = \mu_r \mu_0 \vec{H}$$

$$= (9901)(4\pi \cdot 10^{-7}) 4\vec{a}_z$$

$$= 49.8 \text{ mWb/m}^2 \quad \dots(4)$$

3.19/ Given $\chi_m = 59$,
 $\vec{B} = 0.01 \vec{a}_x$ Wb/m².

$$\mu_r = 1 + \chi_m = 60.$$

From $\vec{B} = \mu\vec{H} = \mu_r \mu_0 \vec{H}$

$$\Rightarrow \vec{H} = \frac{\vec{B}}{\mu_r \mu_0} = \frac{0.01}{60(4\pi \cdot 10^{-7})} \vec{a}_x$$

$$= \vec{a}_x 132.63 \text{ A/m.}$$

THE VOLUME MAGNETIZATION DENSITY

W BE FOUND FROM:

$$\vec{M} = \chi_m \vec{H} = 59(132.63)$$

$$= 7,825 \vec{a}_x \text{ A/m.}$$

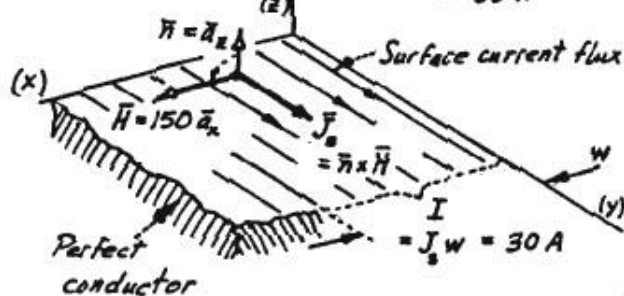
$$\vec{H} = \frac{\vec{B}}{\mu_0} - \vec{M} \Rightarrow \vec{H} = \frac{\vec{B}}{\mu_0} - \vec{M} = 7,825 \vec{a}_x$$

A/m

3-24.) (a) From (3-72).

$$\vec{J}_s = \vec{n} \times \vec{H} = \vec{a}_z \times \vec{a}_x 150 = 150 \vec{a}_y \text{ A/m}$$

In $w = 20$ cm width, $I = J_s w = 150(0.2)$
 $= 30 \text{ A}$

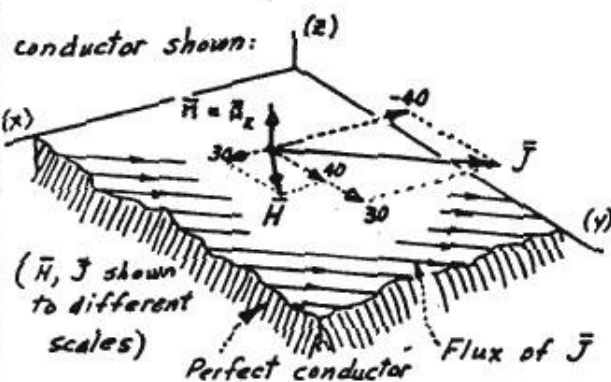


(b) Given $\vec{H} = \vec{a}_x 30 + \vec{a}_y 40$ A/m:

$$\vec{J}_s = \vec{n} \times \vec{H} = \vec{a}_z \times (\vec{a}_x 30 + \vec{a}_y 40)$$

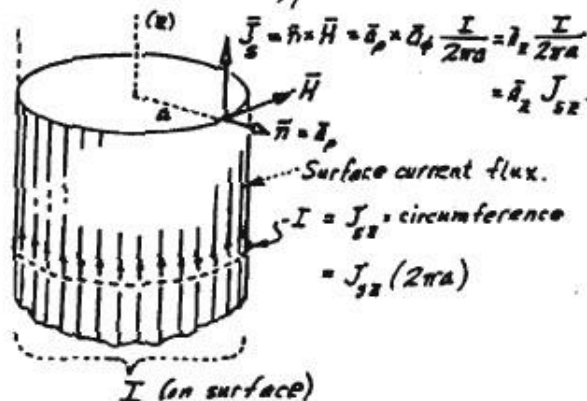
$$= -40 \vec{a}_x + 30 \vec{a}_y \text{ A/m}$$

on the surface $z=0$ of the perfect



(c) Given is $B_\phi = \frac{\mu_0 I}{2\pi r}$, or $\vec{H} = \vec{a}_\phi \frac{I}{2\pi r}$

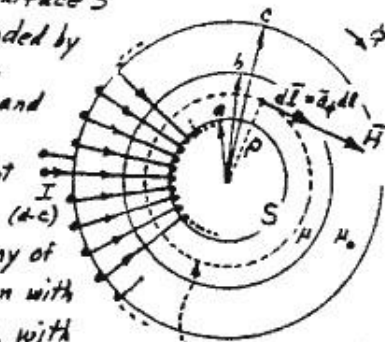
outside a round, perfect conductor:



Problem # 4

3-26.) (The surface S

shown is bounded by the symmetric closed path ℓ , and is pierced by the total current $i = nI$.)



(1/4)

(a) Grasping any of the wires shown with the right hand, with the thumb thus pointing in the direction of I , the fingers within the interior of the toroid then point in the direction of \vec{H} produced by I . \vec{H} is thus ϕ -directed as shown (with the z axis for this circular cylindrical system into the paper).

(b) Ampere's integral law in static form, from (3-66), is applied to the symmetric closed path ℓ shown; so $\oint \vec{H} \cdot d\vec{\ell} = \int_S \vec{J} \cdot d\vec{s} = i$ becomes

$$\oint \vec{H} \cdot d\vec{\ell} = \oint (\vec{a}_\phi H_\phi) \cdot \vec{a}_\phi d\ell = i = nI$$

and with H_ϕ constant on the symmetric ℓ , solving for it obtains

$$H_\phi = \frac{nI}{\oint d\ell} = \frac{nI}{2\pi\rho} \quad \dots(1) \quad (1/2)$$

a result correct for any radius ρ within the toroidal winding whether in the air or the iron region, since no more current is intercepted by ℓ upon expanding its radius into the air region ($b < \rho < c$).

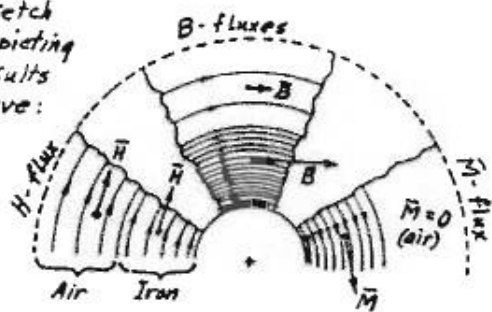
\vec{B} is next found from (3-64c), $\vec{B} = \mu\vec{H}$. Thus,

$$\vec{B} = \begin{cases} \mu\vec{H} = \vec{a}_\phi \frac{\mu nI}{2\pi\rho} & (\text{in Fe: } a < \rho < b) \quad \dots(2) \quad (1/2) \\ \mu_0\vec{H} = \vec{a}_\phi \frac{\mu_0 nI}{2\pi\rho} & (\text{in air: } b < \rho < c) \quad \dots(3) \quad (1/2) \end{cases}$$

(c) The volume magnetization intensity \vec{M} is found from (3-60) and (3-62):

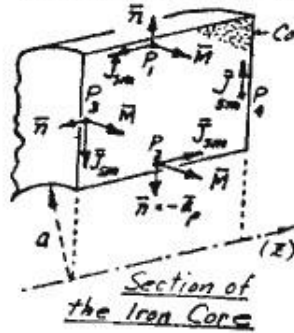
$$\vec{M} = \begin{cases} \chi_m \vec{H} = \vec{a}_\phi (\mu_r - 1) \frac{nI}{2\pi\rho} & (a < \rho < b) \quad \dots(4) \quad (1/4) \\ (\chi_m = 0): = 0 & (\text{air: } b < \rho < c) \quad \dots(5) \quad (1/4) \end{cases}$$

Sketch depicting results above:



(1/2)

(d) With \vec{M} in the iron core $\frac{1}{\rho}$ -dependent, from (4), using (3-56), the volume density of (bound) magnetization current $\vec{J}_m = \nabla \times \vec{M}$ is zero.



From (3-74), the (bound) surface magnetization current densities are given by $\vec{J}_{sm} = -\vec{n} \times \vec{M}$

Thus, at P_2 in the diagram (at $\rho = a$):

$$\vec{J}_{sm} = -\vec{n} \times \vec{M} = \vec{a}_\rho \times \vec{a}_\phi M_\phi = \vec{a}_\phi (\mu_r - 1) \frac{nI}{2\pi a} \quad [A/m] \quad \dots(6) \quad (1/4)$$

Similarly, at $P_1, P_3,$ and P_4 shown:

$$P_1: \vec{J}_{sm} = -\vec{a}_z (\mu_r - 1) \frac{nI}{2\pi b} \quad \dots(7) \quad (1/4)$$

$$P_3: \vec{J}_{sm} = -\vec{a}_\rho (\mu_r - 1) \frac{nI}{2\pi\rho} \quad (a < \rho < b) \quad \dots(8) \quad (1/4)$$

$$P_4: \vec{J}_{sm} = \vec{a}_\rho (\mu_r - 1) \frac{nI}{2\pi\rho} \quad (a < \rho < b) \quad \dots(9) \quad (1/4)$$

(e) At $\rho = a = 1 \text{ cm}$ (in Fe core):

$$\vec{H} = \vec{a}_\phi \frac{nI}{2\pi a} = \vec{a}_\phi \frac{100(0.1)}{2\pi(0.01)} = \vec{a}_\phi 159.2 \text{ A/m} \quad (1/4)$$

$$\vec{B} = \mu\vec{H} = \vec{a}_\phi (10^3) 4\pi \cdot 10^{-7} \cdot 159.2 = \vec{a}_\phi 0.2 \text{ Wb/m}^2 \quad (1/4)$$

At $\rho = b = 1.5 \text{ cm}$ (in Fe core):

$$\vec{H} = \vec{a}_\phi \frac{10}{2\pi(0.015)} = 106.1 \vec{a}_\phi \text{ A/m}, \quad \vec{B} = 0.133 \text{ Wb/m}^2 \quad (1/4)$$

At $\rho = b = 1.5 \text{ cm}$ (in air core):

$$\vec{H} = 106.1 \vec{a}_\phi \text{ A/m}, \quad \vec{B} = \mu_0 \vec{H} = 4\pi \cdot 10^{-7} (106.1) = 0.133 \text{ mWb/m}^2 \quad (1/4)$$

At $\rho = c = 2 \text{ cm}$ (in air core):

$$\vec{H} = \vec{a}_\phi \frac{10}{2\pi(0.02)} = 79.6 \vec{a}_\phi \text{ A/m}, \quad \vec{B} = 0.1 \text{ mWb/m}^2 \quad (1/4)$$

(5)

Problem # 5

3-28) (a) From Ampere's Law AND SYMMETRY:

$$\oint \vec{H} \cdot d\vec{l} = \int \vec{J} \cdot d\vec{s} = i, \quad \vec{H} = \vec{a}_\phi H_\phi$$

For $0 < \rho < a$: $\vec{J} = \vec{a}_z J_z$ with $J_z = \frac{I}{\pi a^2}$ [A/m²]

$$\therefore 2\pi\rho H_\phi = \frac{I}{\pi a^2} \int_0^\rho ds = \frac{I}{\pi a^2} \pi\rho^2 = \frac{I\rho^2}{a^2}$$

$$\Rightarrow H_\phi = \frac{I\rho}{2\pi a^2} \text{ A/m}$$

For $a < \rho < b$: $i = I \Rightarrow H_\phi = \frac{I}{2\pi\rho}$

For $\rho > b$: $i = I \Rightarrow H_\phi = \frac{I}{2\pi\rho}$

SUMMARY:

$$\left\{ \begin{array}{l} 0 < \rho < a: \vec{H} = \vec{a}_\phi \frac{I\rho}{2\pi a^2} \text{ A/m}, \quad \vec{B} = \mu_0 \vec{H} = \vec{a}_\phi \frac{\mu_0 I\rho}{2\pi a^2} \text{ Wb/m}^2 \\ a < \rho < b: \vec{H} = \vec{a}_\phi \frac{I}{2\pi\rho} \text{ A/m}, \quad \vec{B} = \mu \vec{H} = \vec{a}_\phi \frac{\mu I}{2\pi\rho} \text{ Wb/m}^2 \\ \rho > b: \vec{H} = \vec{a}_\phi \frac{I}{2\pi\rho} \text{ A/m}, \quad \vec{B} = \mu_0 \vec{H} = \vec{a}_\phi \frac{\mu_0 I}{2\pi\rho} \text{ Wb/m}^2 \end{array} \right.$$

(b) $\mu_r = 1 + \chi_m \Rightarrow \chi_m = \mu_r - 1$

$$\vec{M} = \chi_m \vec{H} = (\mu_r - 1) \vec{H} = \vec{a}_\phi \underbrace{(\mu_r - 1) \frac{I}{2\pi\rho}}_{\text{IN THE MAGNETIC SLEEVES}} \text{ [A/m]}$$

H_ϕ IS CONTINUOUS AT THE INTERFACES

$B_\phi \neq M_\phi$ ARE DISCONTINUOUS AT THE INTERFACES

(N.B. SEE p.4 FOR SKETCH OF $H_\phi, B_\phi, \& M_\phi$ vs. ρ).

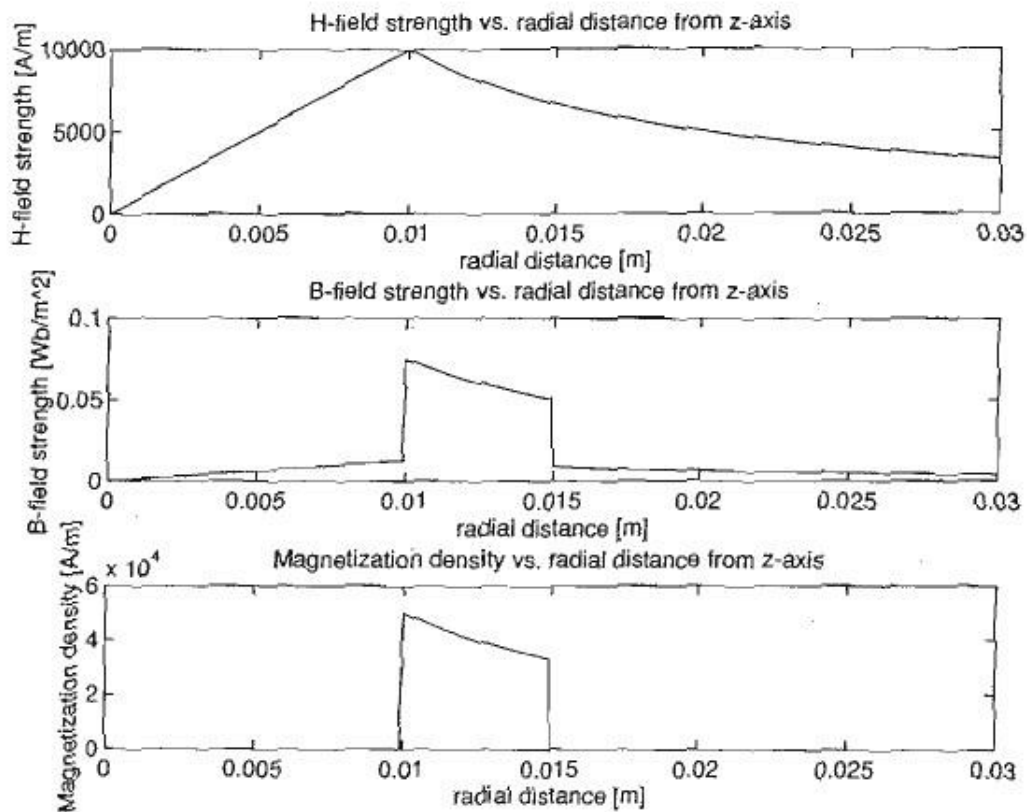
(c) WITHIN THE MAGNETIC SLEEVES:

$$\vec{H} = \vec{a}_\phi \frac{(\mu_r - 1) I}{2\pi\rho}, \quad \therefore \vec{J}_m = \vec{\nabla} \times \vec{H} = 0.$$

ON THE MAGNETIC SLEEVE SURFACE,

AT $\rho = a$: $\vec{J}_{sm} = -n \times \vec{H} = -(-\vec{a}_z) \times \vec{a}_\phi H_\phi = \vec{a}_z (\mu_r - 1) \frac{I}{2\pi\rho}$ [A/m]

AT $\rho = b$: $\vec{J}_{sm} = -n \times \vec{H} = -\vec{a}_\rho \times \vec{a}_\phi H_\phi = -\vec{a}_z (\mu_r - 1) \frac{I}{2\pi\rho}$ [A/m]

Problem # 5 (continued)

$$\#7) \quad \underline{J}_m = \nabla \times \underline{M} = -a_y C$$

i.e. NEGATIVE "y" DIRECTION.

351B ASSIGNMENT #5 - SOLUTIONS OUTLINE

6).

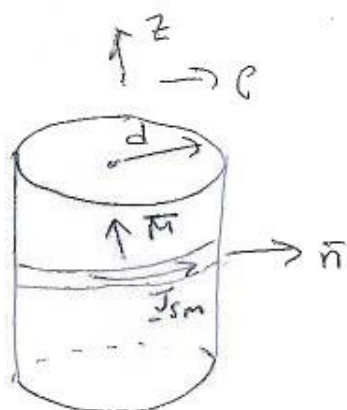


Fig. 1: DIRECTION OF \vec{J}_{sm} FOR UNIFORM MAGNETIZATION

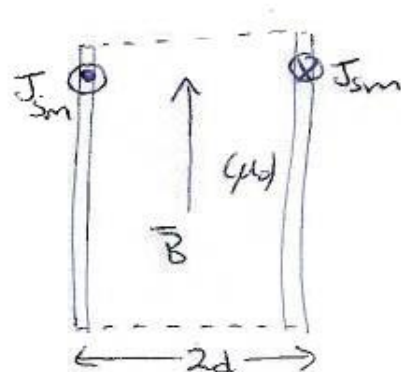


Fig. 2:

EQUIVALENT SURFACE CURRENT MODEL

SINCE \vec{M} IS CONSTANT : $\nabla \times \vec{M} = \vec{J}_m = 0$, BUT THERE IS A BOUND SURFACE MAGNETIZATION CURRENT DENSITY :

$$-\vec{n} \times \vec{M} = \vec{J}_{sm} \quad (\text{A/m})$$

$$\Rightarrow \vec{J}_{sm} = -\vec{a}_\phi \times (5000) \vec{a}_z = \vec{a}_\phi (5000) \quad (\text{A/m})$$

\therefore USING THE EQUIVALENT CURRENT-SHEET MODEL (FIG. 2) & AMPÈRE'S LAW : (SEE P. 45, CHAP. 1. LECTURE NOTES)

$$\vec{B} = 2 \times \frac{\mu_0 J_{sm}}{2} \vec{a}_z \quad \left(\begin{array}{l} \text{INSIDE THE CYLINDER} \\ = 0 \text{ OUTSIDE!} \end{array} \right)$$

$$\vec{B} = 4\pi \times 10^{-7} \times 5000 = 0.00628 \vec{a}_z \quad (\text{T or Wb/m}^2)$$

CONTINUED . . .

6 (... CONTINUED)

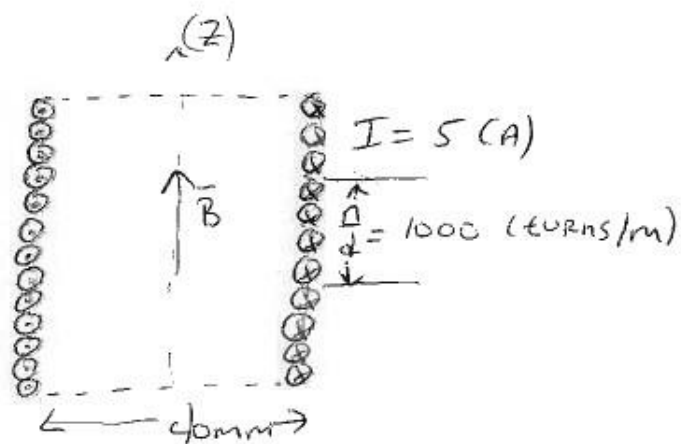


FIG. 3:

AN EQUIVALENT
SOLENOID IMPLEMENTATION
OF THE MODEL IN
FIG. (2).

- THE TOTAL CURRENT / METER LENGTH MUST BE 5000 (A/m), OR 5000 A · turns / meter.
- ∴ ANY CHOICE OF NUMBER OF TURNS IS CORRECT (IN THEORY) IF CHOOSE THE CURRENT APPROPRIATELY.

e.g. CHOOSE 1000 turns / meter

$$\Rightarrow I = 5 \text{ A}$$

(N.B. 1000 turns / meter \Rightarrow 1 mm diameter wire

FOR SOLENOID MADE OF A SINGLE LAYER OF WIRES WOUND TIGHTLY. IS 5 (A) TOO HIGH FOR 1 (mm) diameter wire? $5 \text{ (A)} \Rightarrow 6.4 \times 10^6 \text{ (A/m}^2)$ CURRENT DENSITY, WHICH IS ACCEPTABLE FOR COPPER WIRE!)