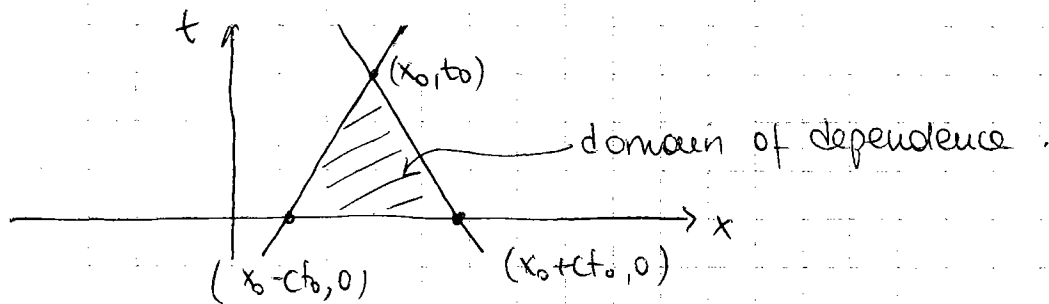
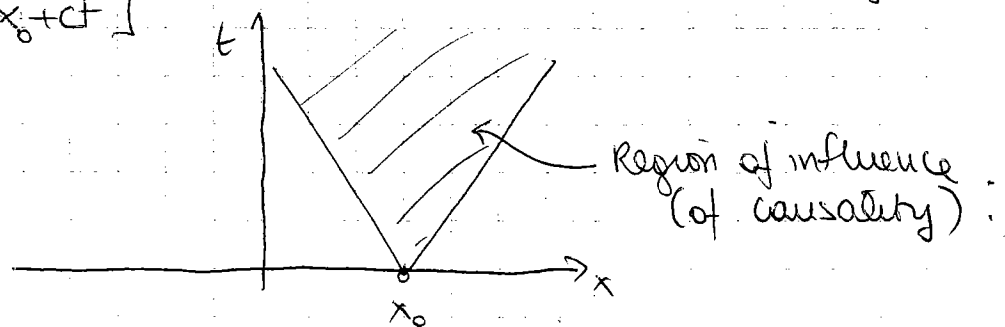


5.3 Interpretation of the characteristics

- Observing the functional form of the solution $u(x,t)$, we see that the value of the solution at a point (x_0, t_0) depends only on the initial conditions within the interval $[x_0 - ct_0, x_0 + ct_0]$.



- Conversely, any point (x_0) in the initial data influences the solution at a later time within a region $[x_0 - ct, x_0 + ct]$.



The characteristics delimit both domain of dependence & region of influence. Information is transported along/within characteristics in a causal way.

Note: if $g(x) = 0 \forall x$ then the characteristics transport all of the information and the solution is an average of the information from $x_0 - ct_0$ and $x_0 + ct_0$.

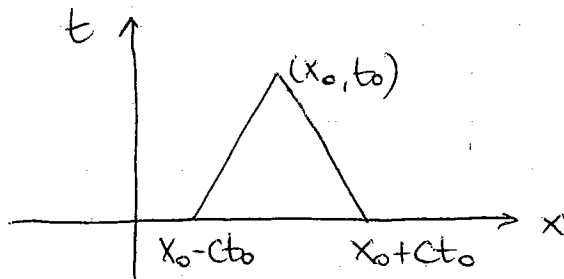
5.4 Non Homogeneous wave equation

(Initial value problem on an ∞ domain)

Suppose we want to solve
$$\begin{cases} u_{tt} - c^2 u_{xx} = F(x, t) \\ u(x, 0) = f(x) \\ u_t(x, 0) = g(x) \end{cases}$$

This corresponds, for instance, to the oscillations of an ∞ string under external forcing.

Let's consider the triangle in the (x, t) plane



and integrate the PDE over the triangle:

$$\iint_{\Delta} (u_{tt} - c^2 u_{xx}) dx dt = \iint_{\Delta} F(x, t) dx dt$$

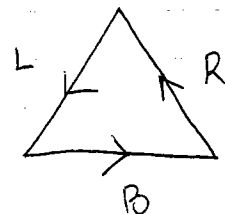
Let $P = u_t$ and $Q = c^2 u_x$ then

$$\iint_{\Delta} (P_t - Q_x) dx dt = \iint_{\Delta} F(x, t) dx dt$$

$$= - \oint_{\Gamma} P(x, t) dx + Q(x, t) dt \quad (\text{following Green's formula})$$

Γ is the contour of the triangle.

Let's define Γ as $B + L + R$
base \uparrow left \uparrow right



Parametric representations of the contour Γ

On the base: $dt = 0$, $t = 0$

$$\begin{aligned} \text{so } \int_B P(x, t) dx + Q(x, t) dt &= \int_B P(x, 0) dx \\ &= \int_B u_t(x, 0) dx = \int_{x_0 - ct_0}^{x_0 + ct_0} g(x) dx \end{aligned}$$

On the left side: $x - ct = \text{constant} = x_0 - ct_0$
so $dx - c dt = 0$

$$\begin{aligned} \text{so } \int_L P(x, t) dx + Q(x, t) dt &= \int_L u_t(x, t) dx + c^2 u_x(x, t) dt \\ &= \int_L u_t(x, t) c dt + c^2 u_x(x, t) \frac{dx}{c} \\ &= c \int_L du = c [-u(x_0, t_0) + u(x_0 - ct_0, 0)] \\ &= c [u(x_0, t_0) + f(x_0 - ct_0)] \end{aligned}$$

not surprising:
L is a characteristic

Similarly on the right side

$$\begin{aligned} \int_R P dx + Q dt &= -c [u(x_0, t_0) - u(x_0 + ct_0, 0)] \\ &= -c [u(x_0, t_0) - f(x_0 + ct_0)] \end{aligned}$$

So, putting it all together we get

$$-2c u(x_0, t_0) + c (f(x_0 - ct_0) + f(x_0 + ct_0)) + \int_{x_0 - ct_0}^{x_0 + ct_0} g(x) dx$$

$$= - \iint_{\Delta} F(x, t) dx dt$$

$$\Rightarrow u(x_0, t_0) = \frac{1}{2} [f(x_0 - ct_0) + f(x_0 + ct_0)] + \frac{1}{2c} \int_{x_0 - ct_0}^{x_0 + ct_0} g(x) dx + \frac{1}{2c} \iint_{\Delta(x, t)} F(x, t) dx dt$$

This is also known as d'Alembert's formula.

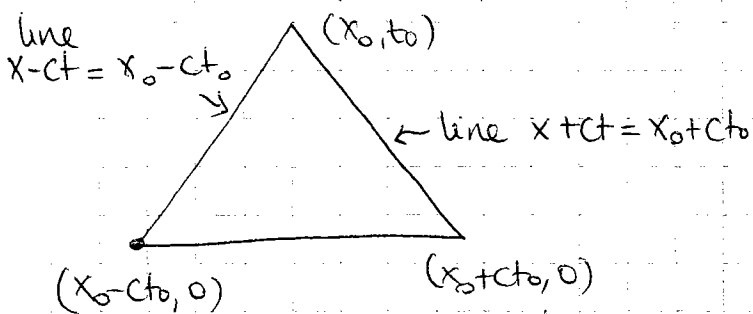
Example:
$$\begin{cases} u_{tt} - 9u_{xx} = 2 \sinh(x) \\ u(x, 0) = x \\ u_t(x, 0) = \sin(x) \end{cases}$$

Here, we use the formula with

$F(x, t) = 2 \sinh(x)$, $c = 3$, $f(x) = x$, $g(x) = \sin x$ so

$$u(x_0, t_0) = \frac{1}{2} [(x_0 - 3t_0) + (x_0 + 3t_0)] + \frac{1}{6} \int_{x_0 - 3t_0}^{x_0 + 3t_0} \sinh(x) dx - \frac{1}{6} \iint_{\Delta} F(x, t) dx dt$$

To perform the double integral, we need to find the bounds.



There are two ways of doing integral: first integrate x then t , or vice versa.

Way # ①:
$$\int_{t=0}^{t=t_0} \int_{x=-ct+x_0+ct_0}^{x=ct+x_0-ct_0} F(x, t) dx dt$$

Way # ②:
$$\int_{x=x_0-ct_0}^{x=x_0} \int_{t=0}^{t=\frac{x-x_0+ct_0}{c}} F(x, t) dx dt + \int_{x=x_0}^{x=x_0+ct_0} \int_{t=0}^{t=\frac{-x+x_0+ct_0}{c}} F(x, t) dx dt$$

Here let's choose #1

$$\begin{aligned} \iint_{\Delta} F &= \int_{t=0}^{t=t_0} 2 [\cosh(-ct + x_0 + ct_0) - \cosh(ct + x_0 - ct_0)] dt \quad \text{with } c=3 \\ &= -\frac{2}{3} \sinh(x_0) - \frac{2}{3} \sinh(x_0) + \frac{2}{3} \sinh(x_0 + 3t_0) + \frac{2}{3} \sinh(x_0 - 3t_0) \end{aligned}$$

$$\begin{aligned} \Rightarrow u(x_0, t_0) &= x_0 + \frac{1}{6} (-\cos(x_0 + 3t_0) + \cos(x_0 - 3t_0)) - \frac{2}{9} \sinh(x_0) \\ &\quad + \frac{1}{9} \sinh(x_0 + 3t_0) + \frac{1}{9} \sinh(x_0 - 3t_0) = x_0 + \frac{1}{3} \sin(x_0) \sin(3t_0) \\ &\quad - \frac{2}{9} \sinh(x_0) + \frac{2}{9} \sinh(x_0) \cosh(3t_0) \end{aligned}$$

5.5 The Energy Method: Existence & Uniqueness of Solutions

5.5.1 Equivalent condition to uniqueness

Consider two associated problems

$$\begin{aligned} \textcircled{P1} : \quad & u_{tt} = c^2 u_{xx} \\ & u(x, 0) = f(x) \quad u(a, t) = h(t) \\ & u_t(x, 0) = g(x) \quad u(b, t) = H(t) \end{aligned}$$

$$\begin{aligned} \textcircled{P2} \quad & u_{tt} = c^2 u_{xx} \\ & u(x, 0) = 0 \quad u(a, t) = 0 \\ & u_t(x, 0) = 0 \quad u(b, t) = 0 \end{aligned}$$

P_2 is said to have homogeneous ICs and BCs.

Lemma: The solution to P_1 is unique
 \Leftrightarrow The only solution to P_2 is $u \equiv 0$.

Proof \Leftarrow (by negation).

Suppose we know two solutions to P_1 , namely $u_1(x, t)$ and $u_2(x, t)$ then

$v(x, t) = u_1(x, t) - u_2(x, t)$ is a non-zero solution of P_2 . \square

The proof of \Rightarrow is similar.

Therefore, to prove that the solution to a given linear PDE is unique, it is enough to show that the solution to the associated problem with homogeneous BCs and ICs is necessarily $u \equiv 0$.

The energy method is an easy way of doing just that.

5.5.2 Example of the Energy Method

Given the PDE $u_{tt} = c^2 u_{xx}$, we now define the function

$$E(t) = \int_a^b [u_t^2 + c^2 u_x^2] dx$$

where a and b are the boundaries of the domain considered.

Then ① $E(t) \geq 0$ (by construction) and

$$\begin{aligned} \text{② } \frac{dE}{dt} &= \int_a^b \frac{\partial}{\partial t} (u_t^2) + c^2 \frac{\partial}{\partial t} (u_x^2) dx \\ &= \int_a^b 2u_t u_{tt} + 2c^2 u_x u_{xt} dx \\ &= \int_a^b [2u_t u_{tt} + 2c^2 \frac{\partial}{\partial x} (u_x u_t) - 2c^2 u_t u_{xx}] dx \\ &= \int_a^b 2c^2 \frac{\partial}{\partial x} (u_x u_t) dx \\ &= 2c^2 [u_x(b,t)u_t(b,t) - u_x(a,t)u_t(a,t)] \end{aligned}$$

If the BCs are such that either u_x or $u_t = 0$ on the boundaries then $\frac{dE}{dt} = 0$ at all times t .

Now, at $t=0$, the "Energy" E of problem P_2 is 0. Since $\frac{dE}{dt} = 0$, then $E=0$ for all times

$$\Rightarrow \text{for all } t, \int_a^b u_t^2 + c^2 u_x^2 dx = 0$$

$$\Rightarrow u_t = u_x = 0 \text{ for all } t \text{ and } x$$

$$\Rightarrow u = 0 \text{ for all } x \text{ and } t.$$

Note: This general concept can also be applied to other PDEs. The idea is derived from the physical concept of energy, but for most PDEs $E(t)$ only needs to satisfy

$$\begin{cases} E(t) \geq 0 \\ \frac{dE}{dt} \leq 0 \\ E(t) \text{ depends on } u(x,t) \text{ and/or its} \\ \text{derivatives.} \end{cases}$$

For the heat equation, for example,

$$\begin{cases} u_t = k u_{xx} \\ u(0,t) = u(L,t) = 0 \end{cases}$$

$$\rightarrow \text{we use } E(t) = \frac{1}{2} \int_0^L u^2(x,t) dx.$$

\rightarrow homework: Prove that it has the required properties